

Nordic Sound Symposium XXIII  
Bolkesjø, Norway  
Sept. 27 - Sept. 30, 2007

# IMPULSE RESPONSE MEASUREMENTS

**Angelo Farina**

Industrial Engineering Dept.  
University of Parma - ITALY  
[HTTP://www.angelifarina.it](http://www.angelifarina.it)



University  
of Parma



# Time Line

## The Past

- Traditional time-domain measurements with pulsive sounds and omnidirectional transducers
- Electroacoustical measurements employing special computer-based hardware, a loudspeaker and an omnidirectional microphone

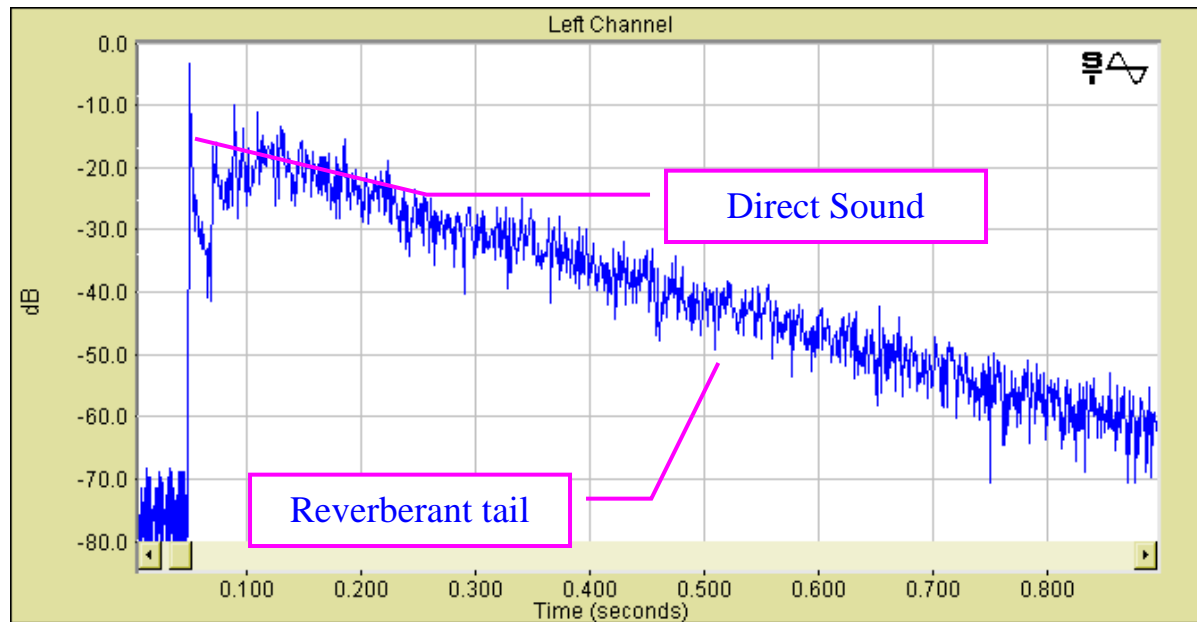
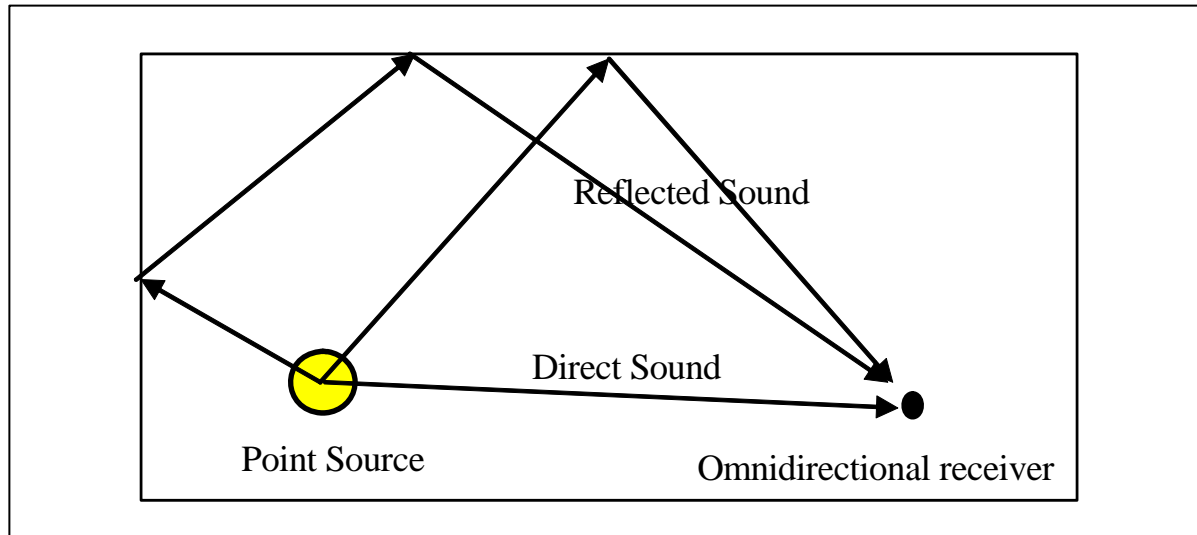
## The Present

- Electroacoustical measurements employing standard sound cards, 2 or more loudspeakers and multiple microphones (2 to 8)

## The Future

- Microphone arrays for capturing high-order spatial information
- Artificial sound sources employing a dense array of loudspeakers, capable of synthesizing the directivity pattern of any real-world source

# Basic sound propagation scheme





# The Past

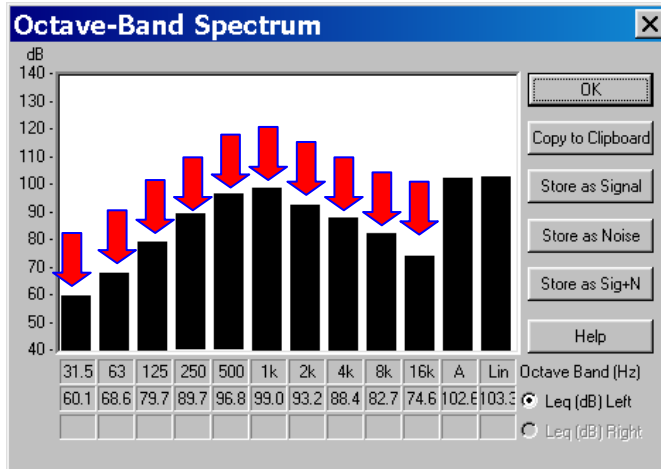
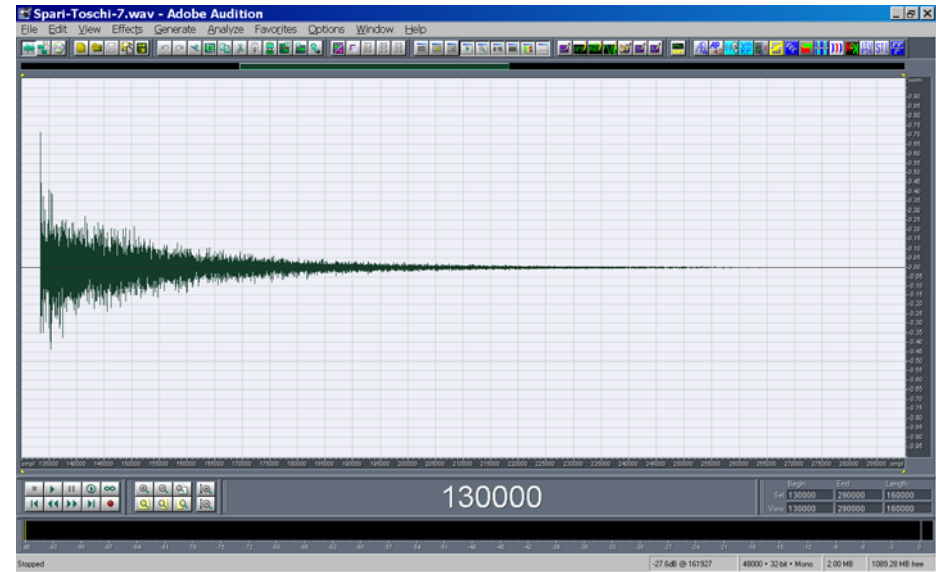
# Traditional measurement methods



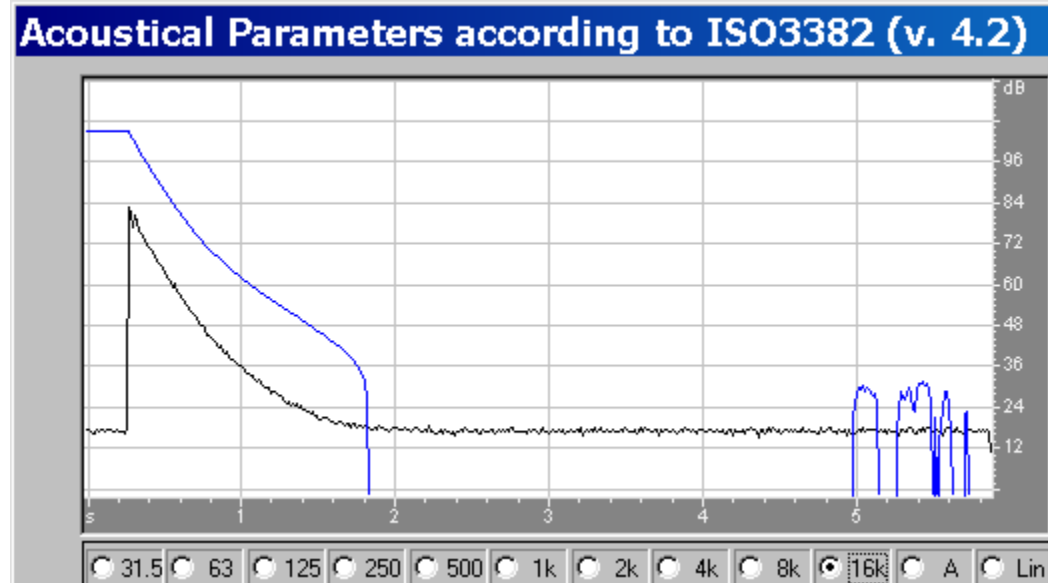
- Pulsive sources: balloons, blank pistol



# Example of a pulsive impulse response



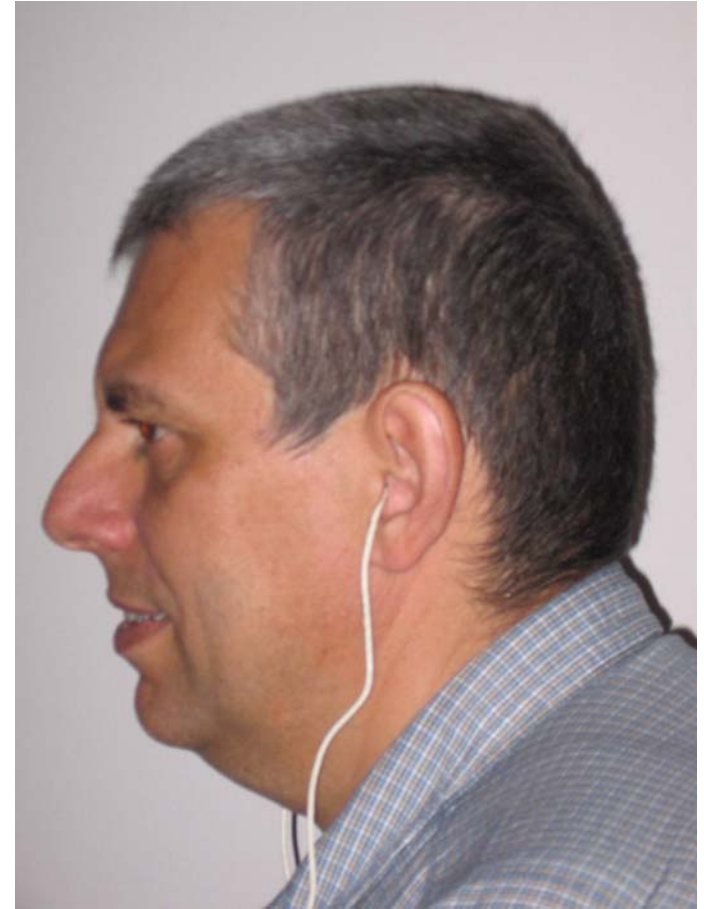
Note the bell-shaped spectrum







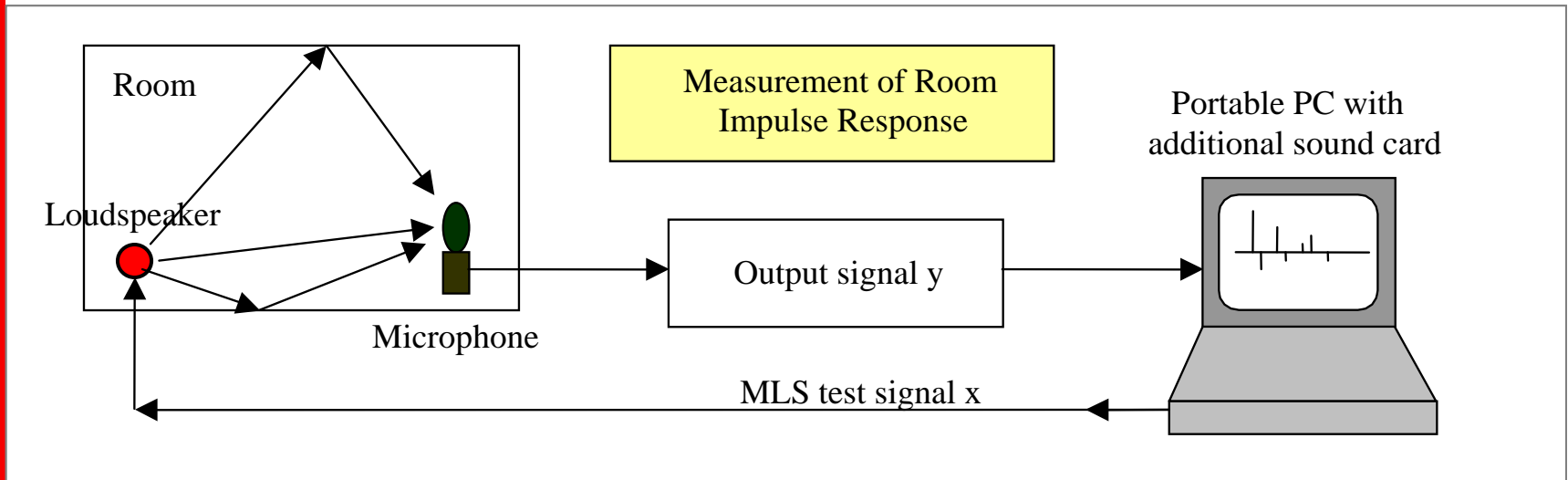
# Test with binaural microphones



- **Cheap electret mikes in the ear ducts**



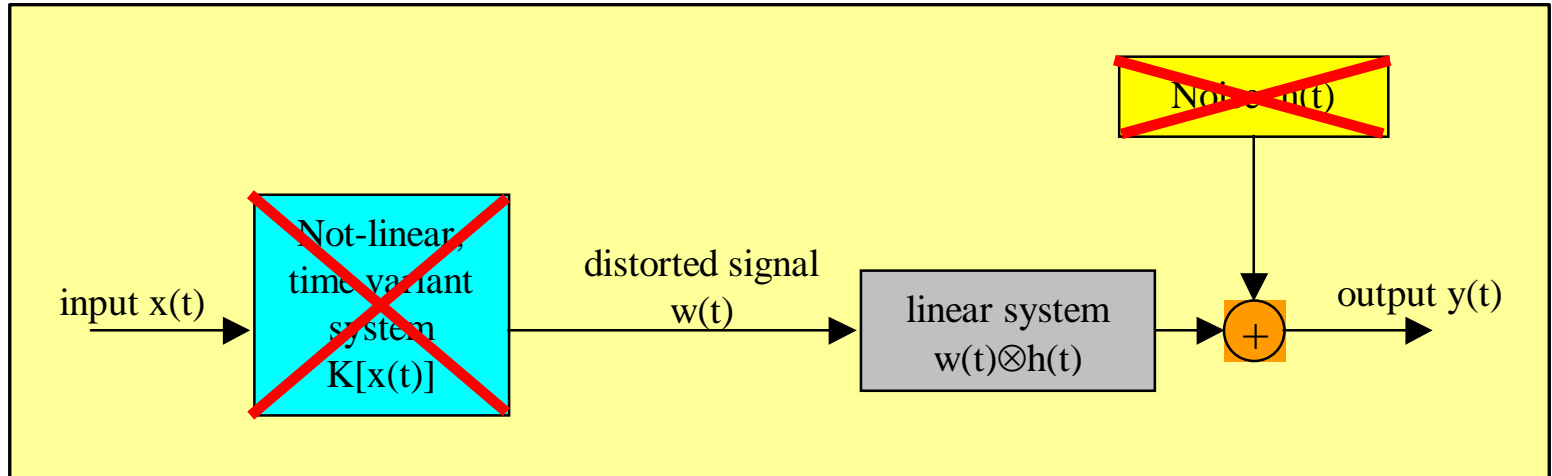
# Loudspeaker as sound source



- A loudspeaker is fed with a special test signal  $x(t)$ , while a microphone records the room response
- A proper deconvolution technique is required for retrieving the impulse response  $h(t)$  from the recorded signal  $y(t)$



# Measurement process



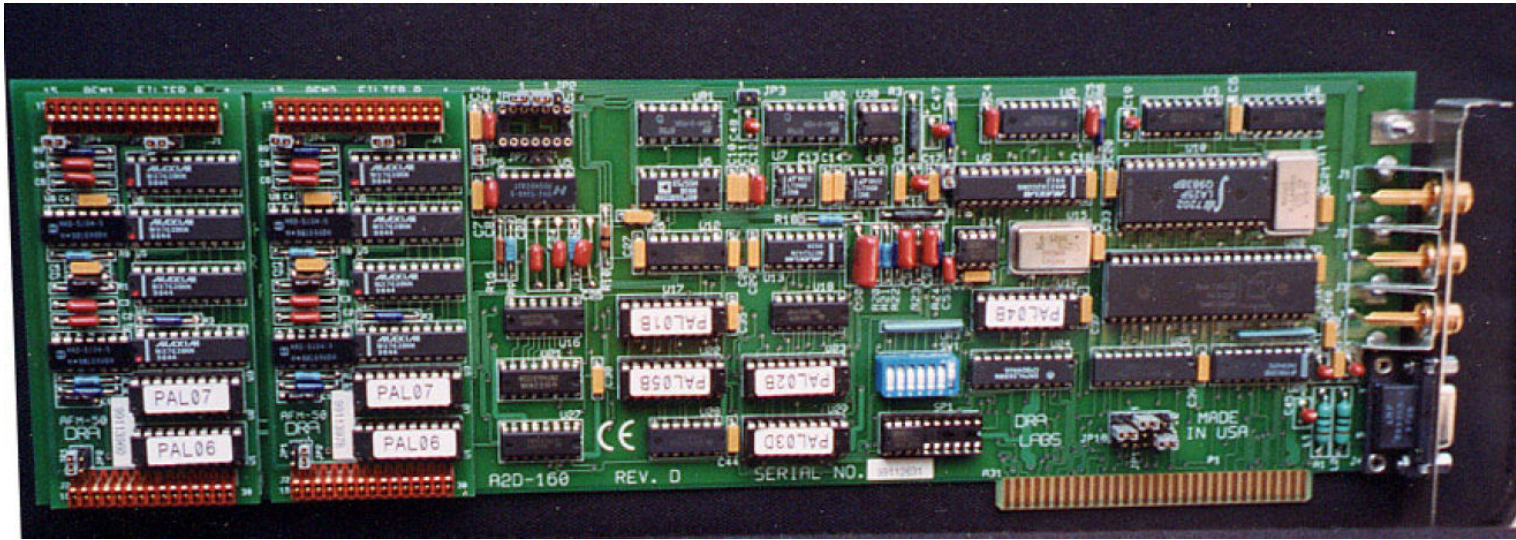
- The desired result is the linear impulse response of the acoustic propagation  $h(t)$ . It can be recovered by knowing the test signal  $x(t)$  and the measured system output  $y(t)$ .
- It is necessary to exclude the effect of the not-linear part  $K$  and of the background noise  $n(t)$ .



# Electroacoustical methods

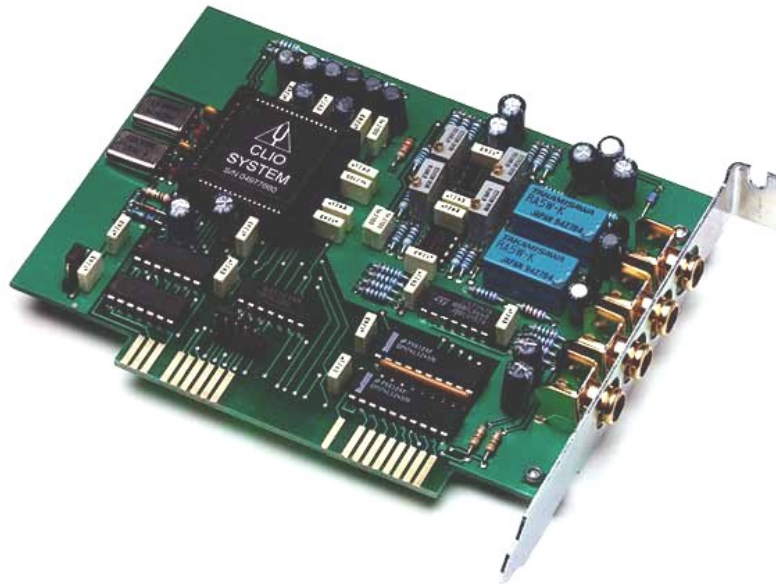
- **Different types of test signals have been developed, providing good immunity to background noise and easy deconvolution of the impulse response:**
  - ▶ MLS (Maximum Length Sequence, pseudo-random white noise)
  - ▶ TDS (Time Delay Spectrometry, which basically is simply a linear sine sweep, also known in Japan as “stretched pulse” and in Europe as “chirp”)
  - ▶ ESS (Exponential Sine Sweep)
- **Each of these test signals can be employed with different deconvolution techniques, resulting in a number of “different” measurement methods**
- **Due to theoretical and practical considerations, the preference is nowadays generally oriented for the usage of ESS with not-circular deconvolution**

# The first MLS apparatus - MLSSA



- MLSSA was the first apparatus for measuring impulse responses with MLS

# More recently - the CLIO system



PB-4281



SC-01

- The Italian-made CLIO system has superseded MLSSA for most electroacoustics applications (measurement of loudspeakers, quality control)



# The first TDS apparatus - TEF



- Techron TEF 10 was the first apparatus for measuring impulse responses with TDS
- Subsequent versions (TEF 20, TEF 25) also support MLS



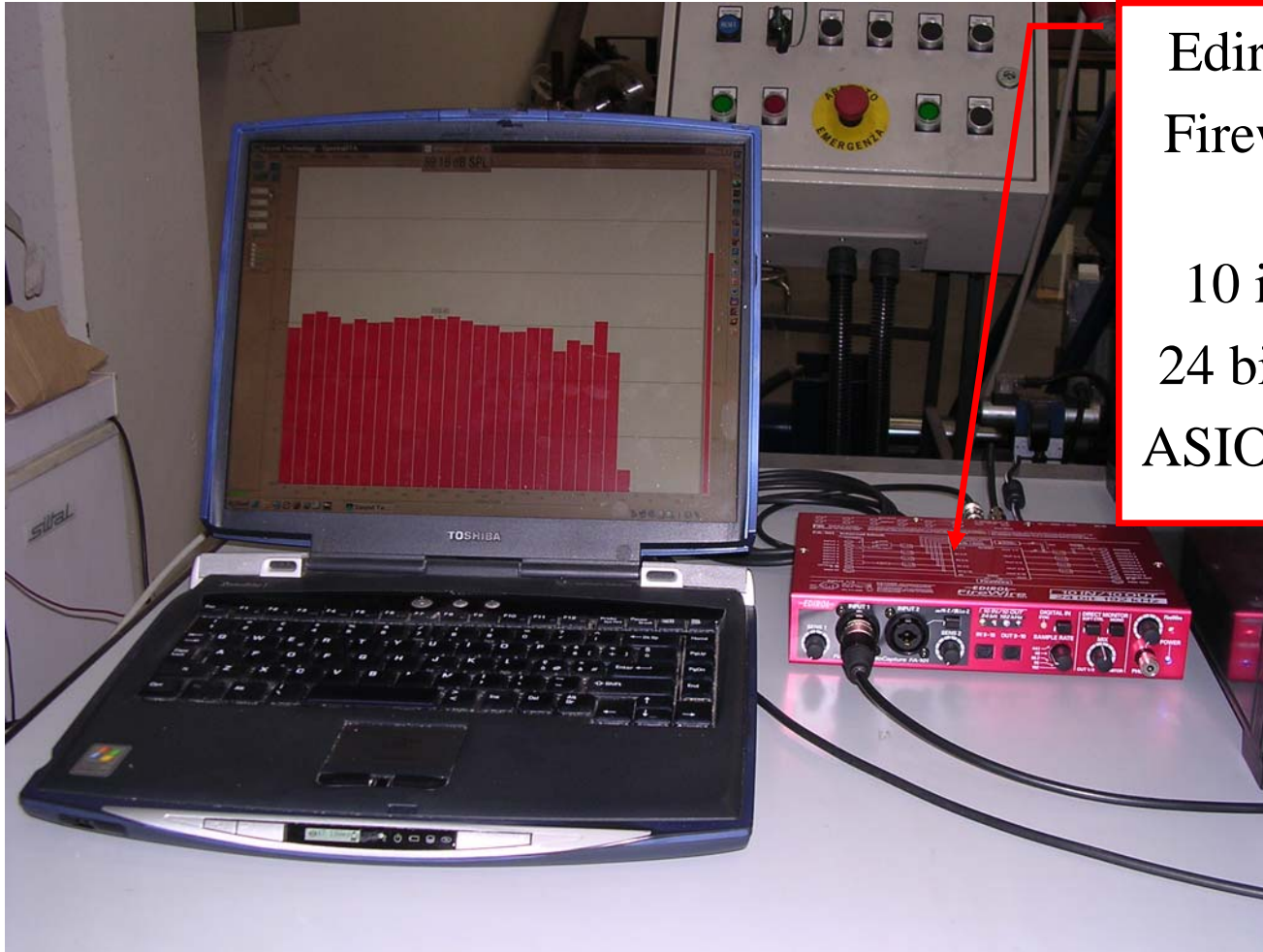


# The Present





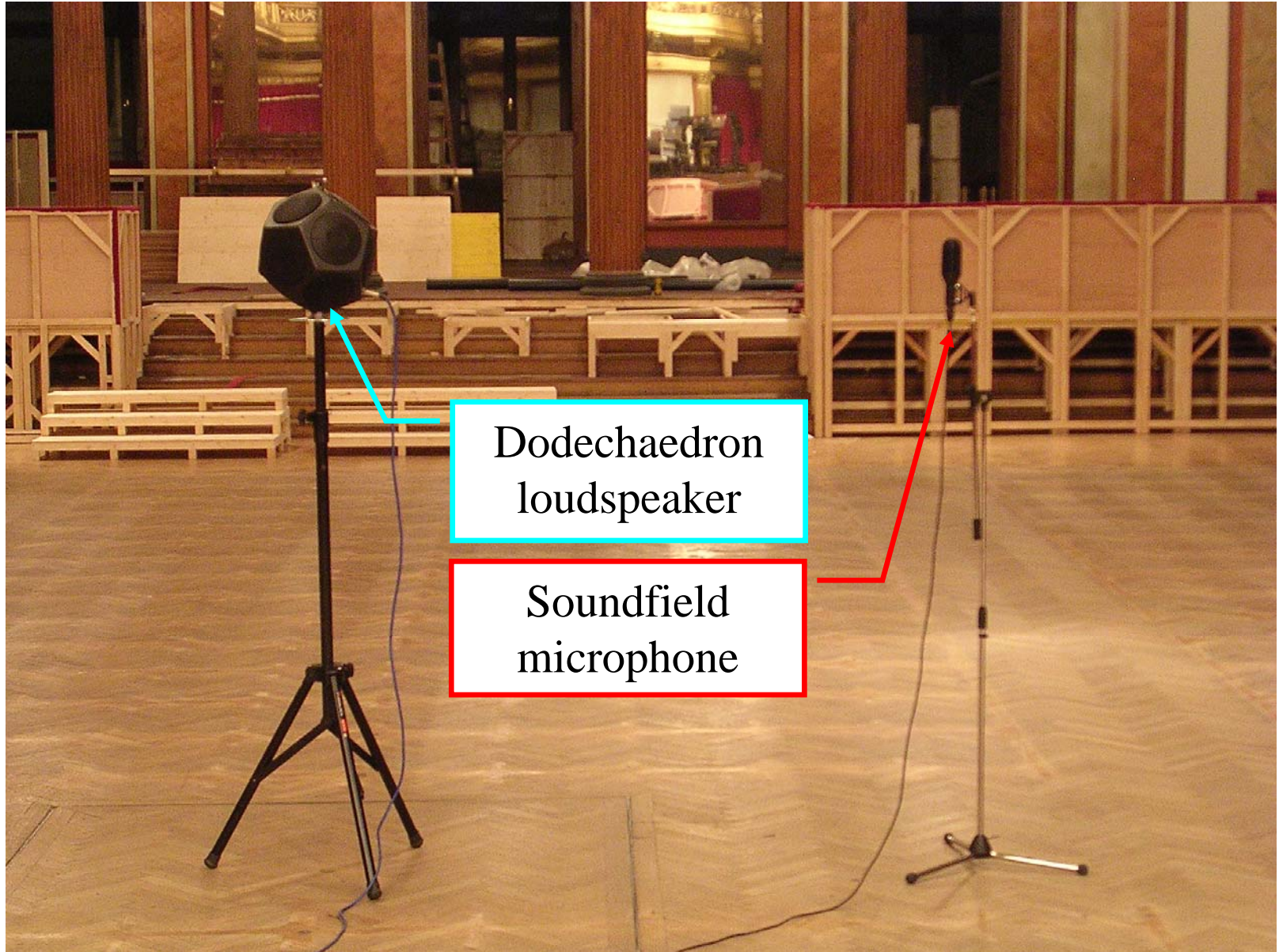
# Today's Hardware: PC and audio interface



Edirol FA-101  
Firewire sound  
card:  
10 in / 10 out  
24 bit, 192 kHz  
ASIO and WDM



# Hardware: loudspeaker & microphone

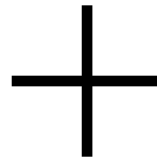
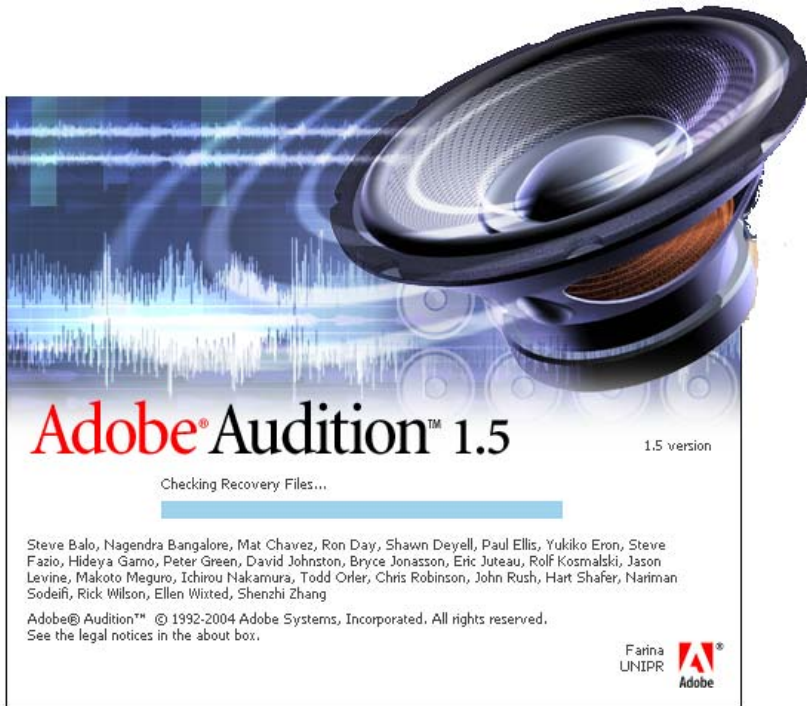






# The first ESS system - AURORA

## Aurora Plugins

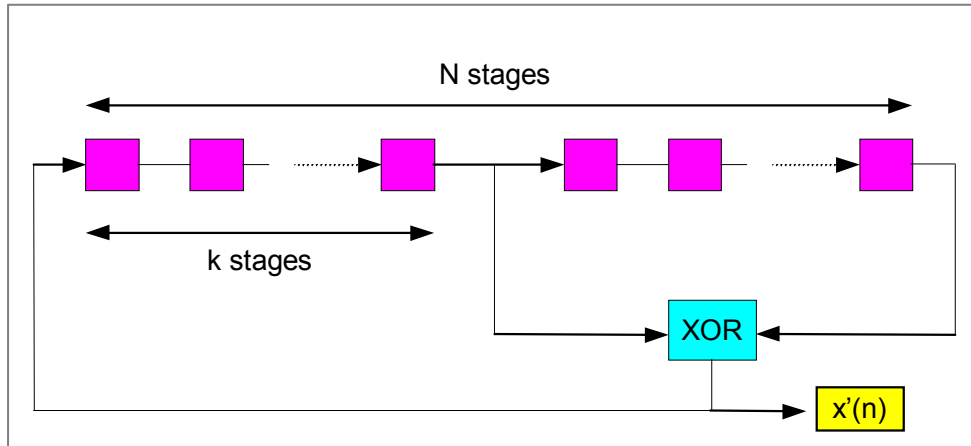


Generate MLS	
Deconvolve MLS	
Generate Sweep	
Deconvolve Sweep	
Convolution	
Kirkeby Inverse Filter	
Speech Transm. Index	

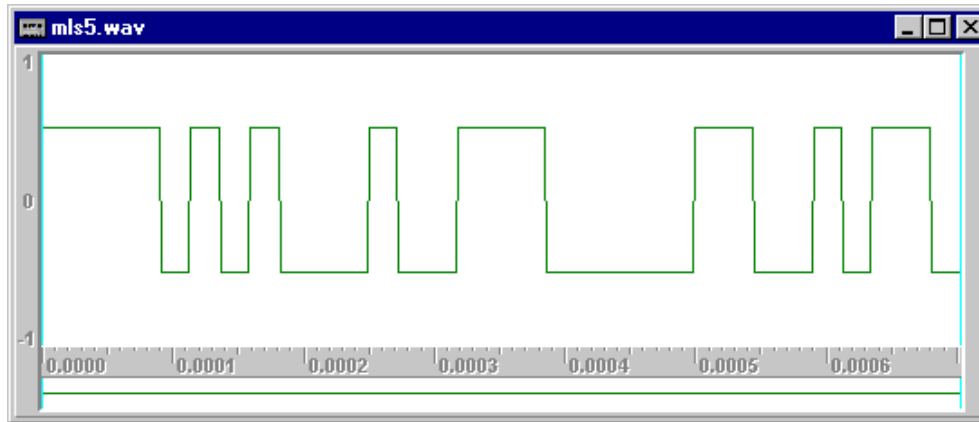
- **Aurora was the first measurement system based on standard sound cards and employing the Exponential Sine Sweep method (1998)**
- **It also works with traditional TDS and MLS methods, so the comparison can be made employing exactly the same hardware**



# MLS method



- **$X(t)$  is a periodic binary signal obtained with a suitable shift-register, configured for maximum length of the period.**



$$L = 2^N - 1$$



- The re-recorded signal  $y(i)$  is cross-correlated with the excitation signal thanks to a fast Hadamard transform. The result is the required impulse response  $h(i)$ , if the system was linear and time-invariant

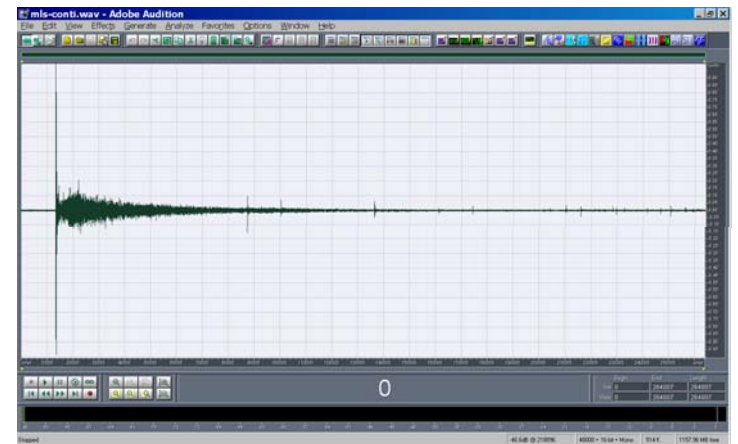
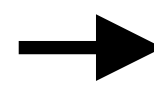
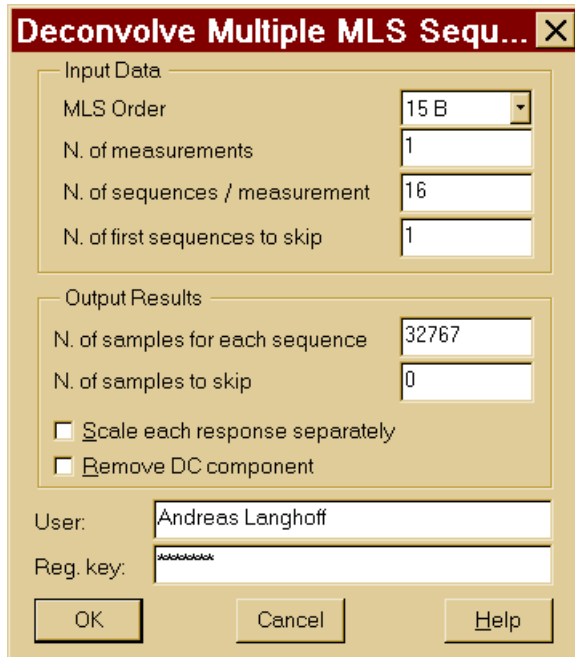
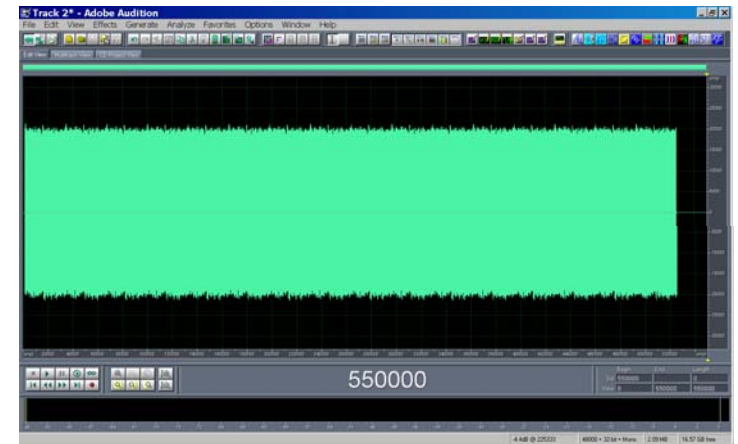
$$h = \frac{1}{L+1} \cdot \tilde{M} \cdot y$$

- Where  $M$  is the Hadamard matrix, obtained by permutation of the original MLS sequence  $m(i)$

$$\tilde{M}(i, j) = m[(i + j - 2) \bmod L] - 1$$



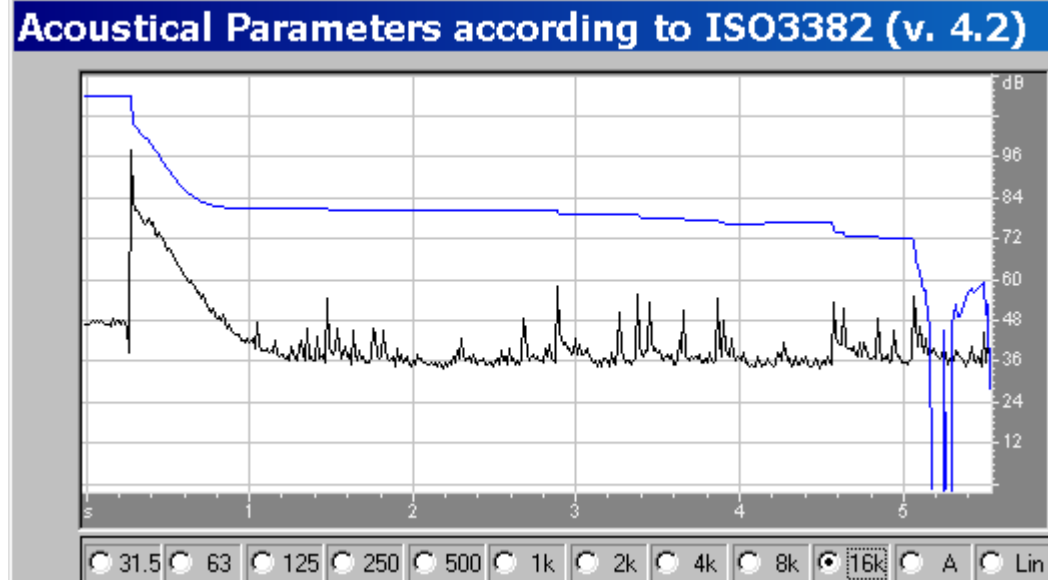
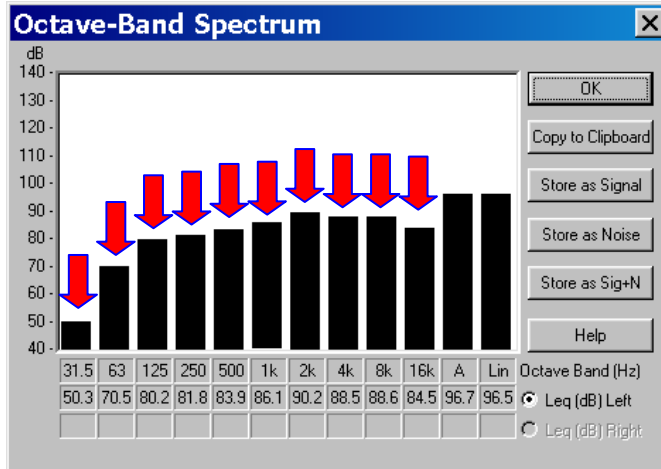
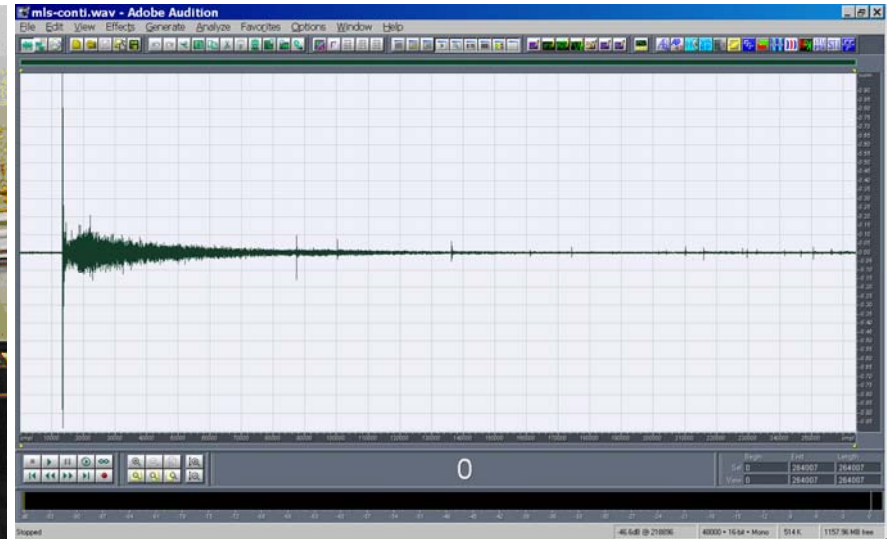
# MLS measurement procedure







# Example of a MLS impulse response





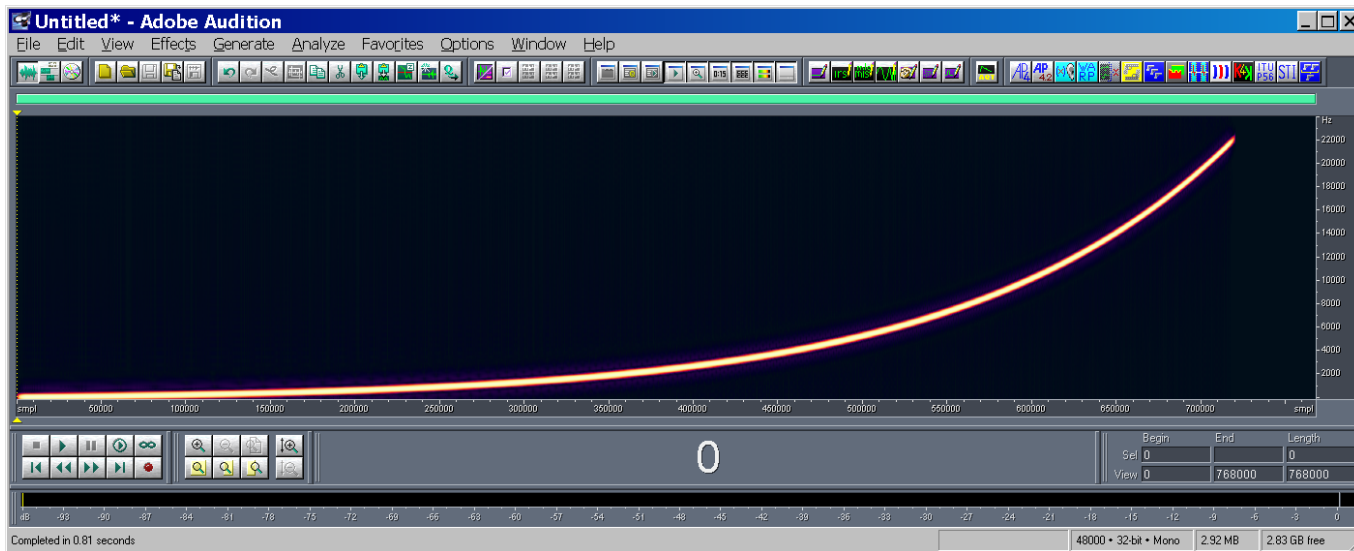
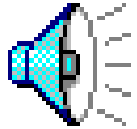
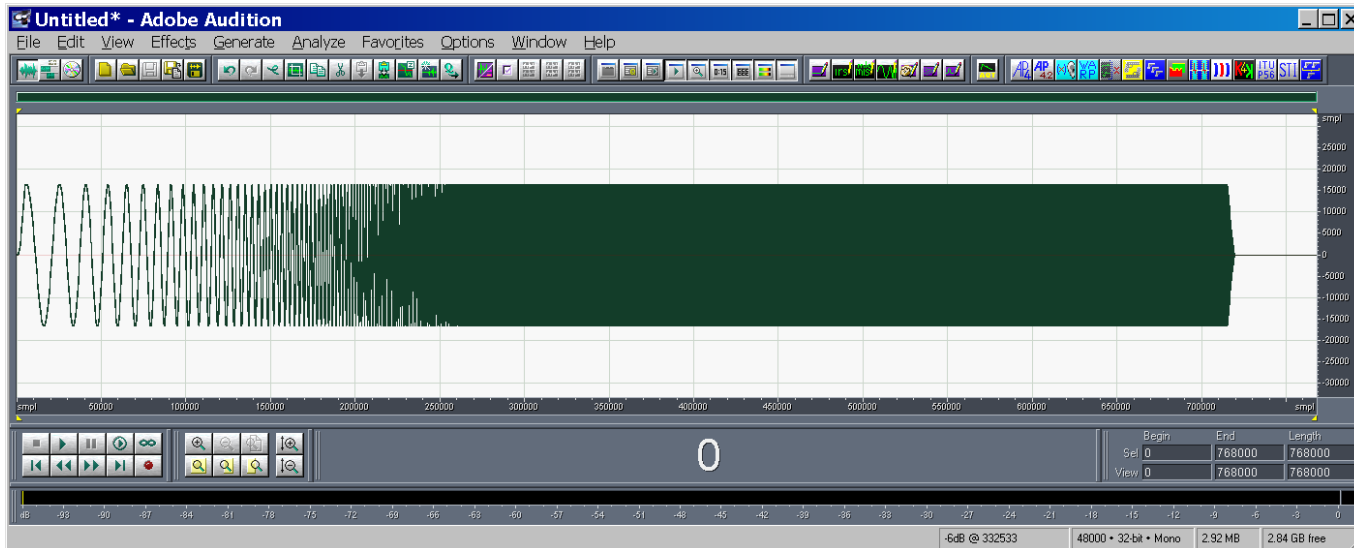
# Exponential Sine Sweep method

- **$x(t)$  is a band-limited sinusoidal sweep signal, which frequency is varied exponentially with time, starting at  $f_1$  and ending at  $f_2$ .**

$$x(t) = \sin \left[ \frac{2 \cdot \pi \cdot f_1 \cdot T}{\ln \left( \frac{f_2}{f_1} \right)} \cdot \left( e^{\frac{t}{T} \cdot \ln \left( \frac{f_2}{f_1} \right)} - 1 \right) \right]$$

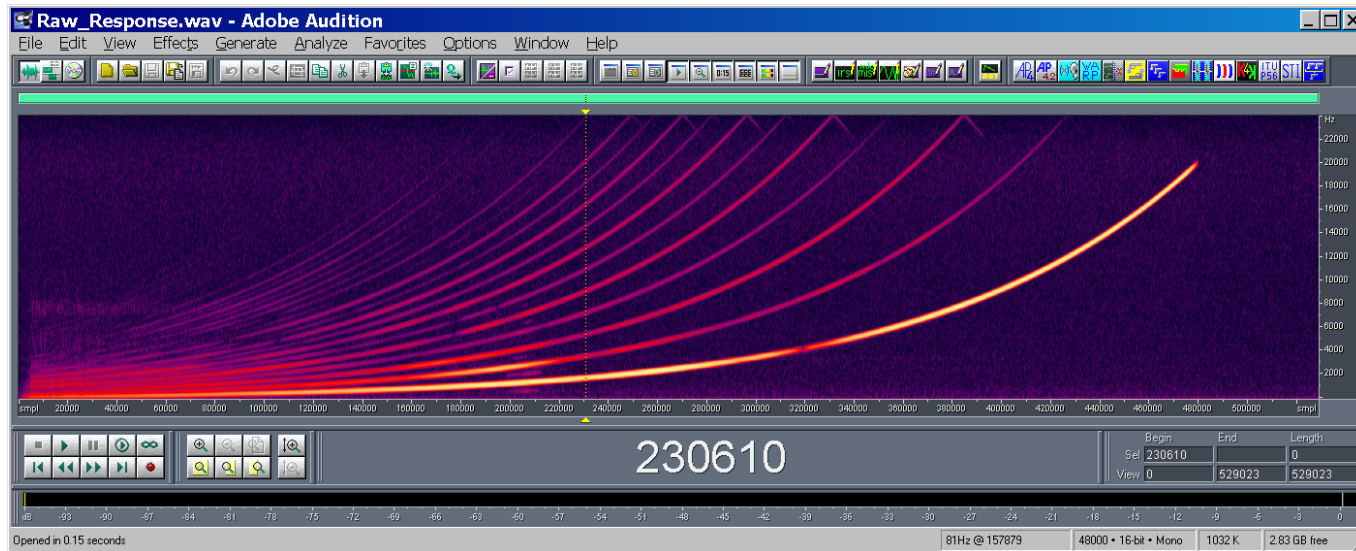
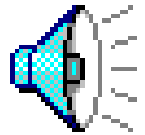
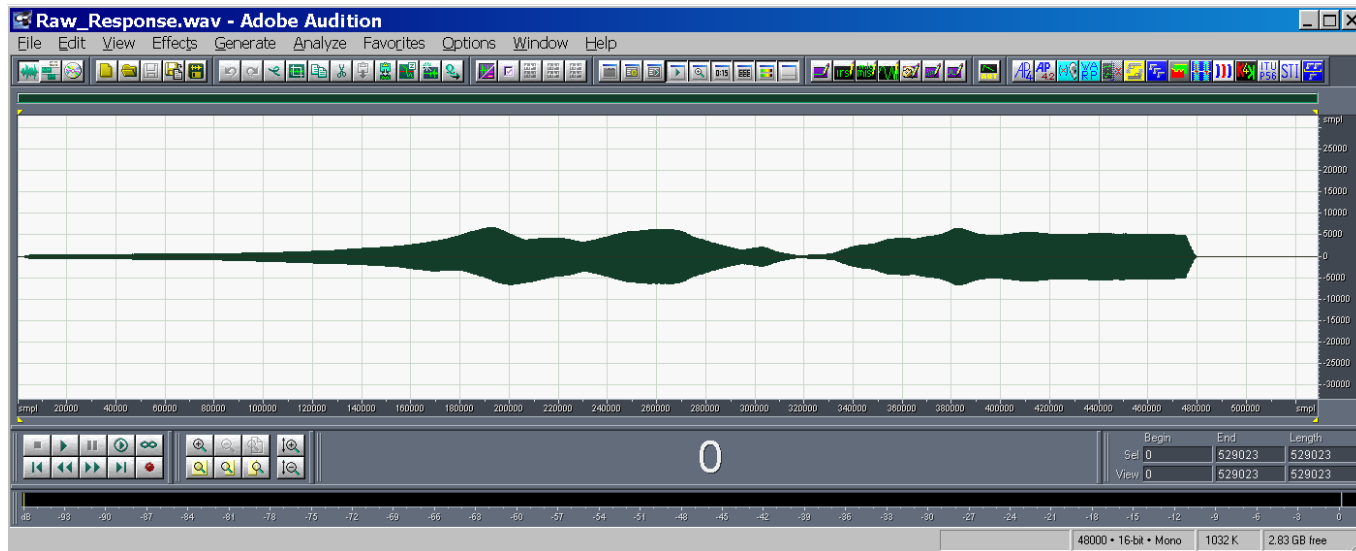


# Test Signal – $x(t)$





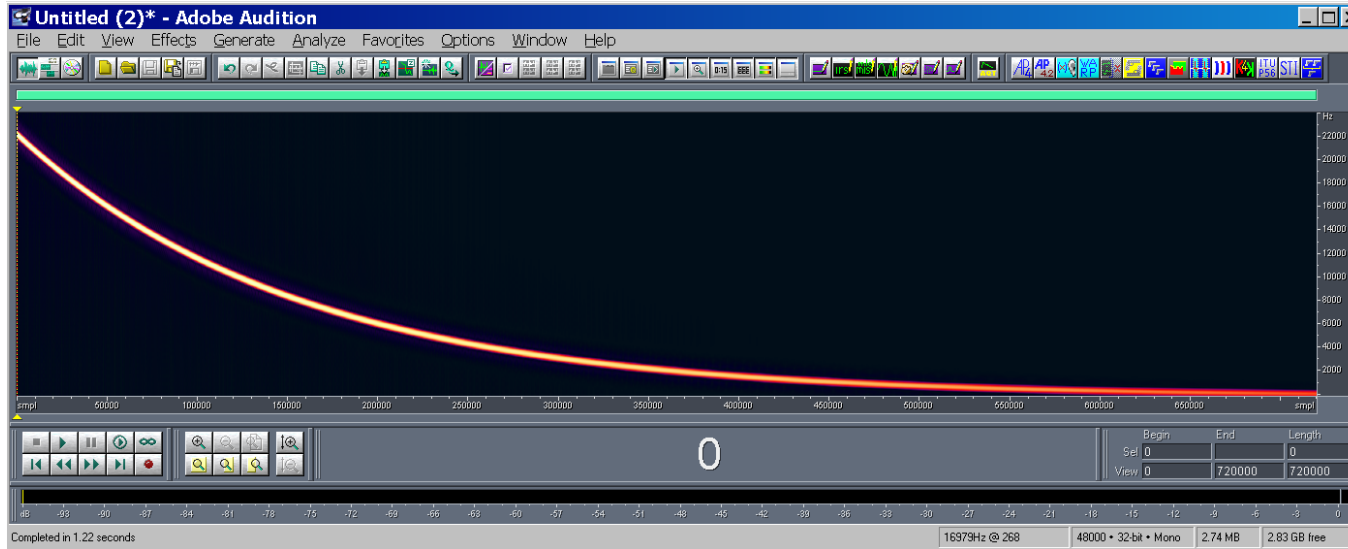
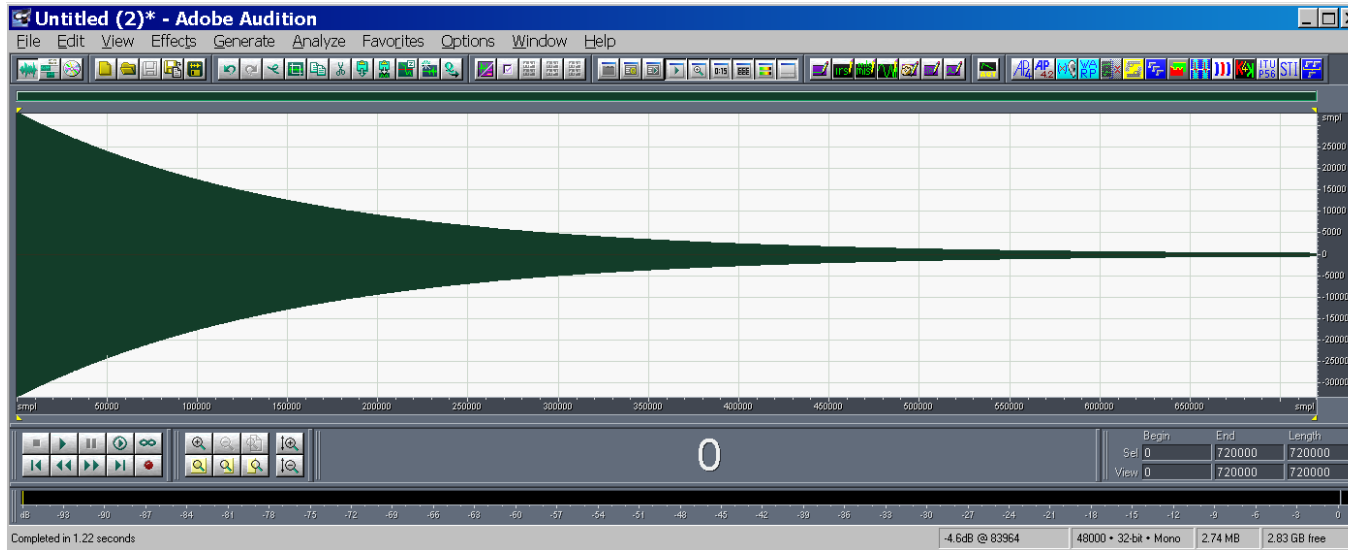
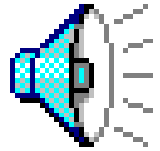
# Measured signal - $y(t)$



- The not-linear behaviour of the loudspeaker causes many harmonics to appear



# Inverse Filter – $z(t)$

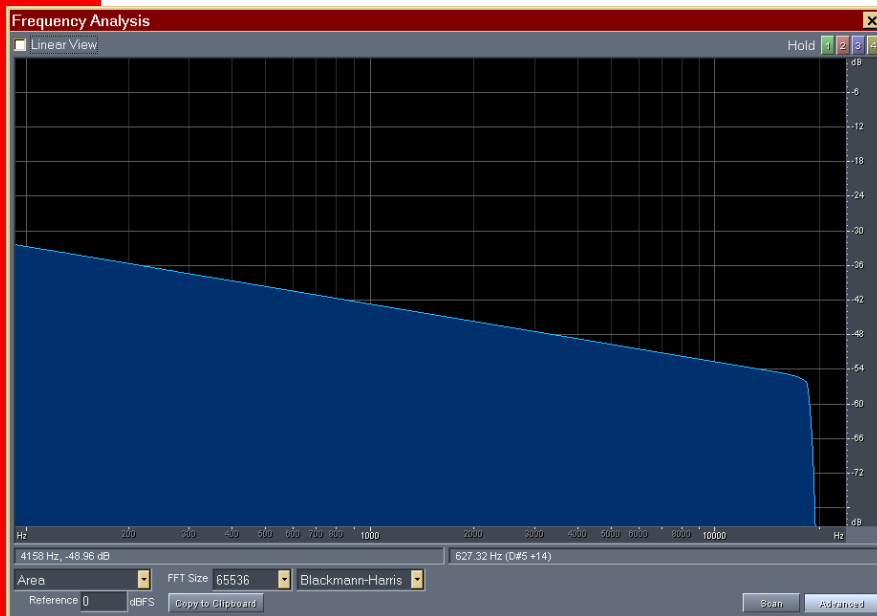


**The deconvolution of the IR is obtained convolving the measured signal  $y(t)$  with the inverse filter  $z(t)$  [equalized, time-reversed  $x(t)$ ]**

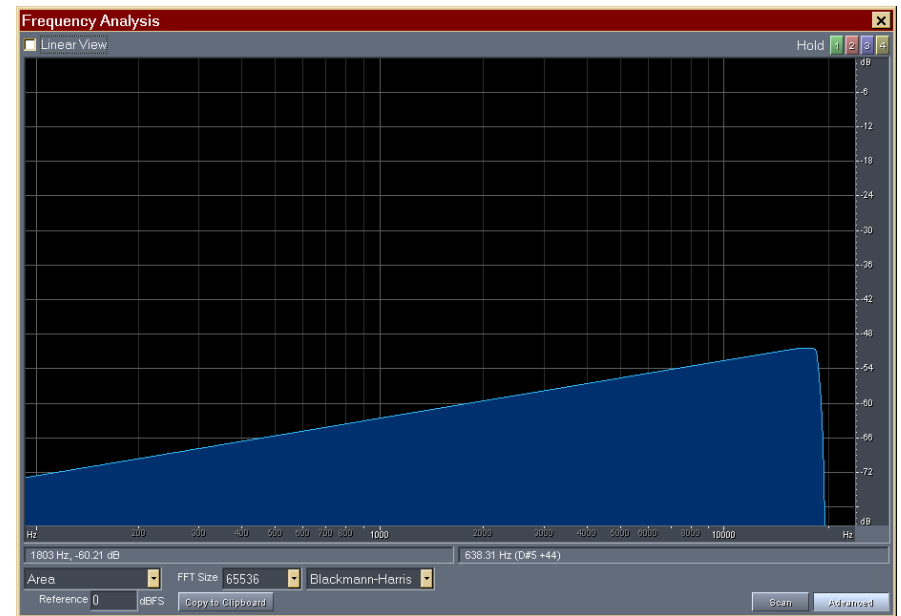
# Deconvolution of Exponential Sine Sweep



The “time reversal mirror” technique is employed: the system’s impulse response is obtained by convolving the measured signal  $y(t)$  with the time-reversal of the test signal  $x(-t)$ . As the log sine sweep does not have a “white” spectrum, proper equalization is required



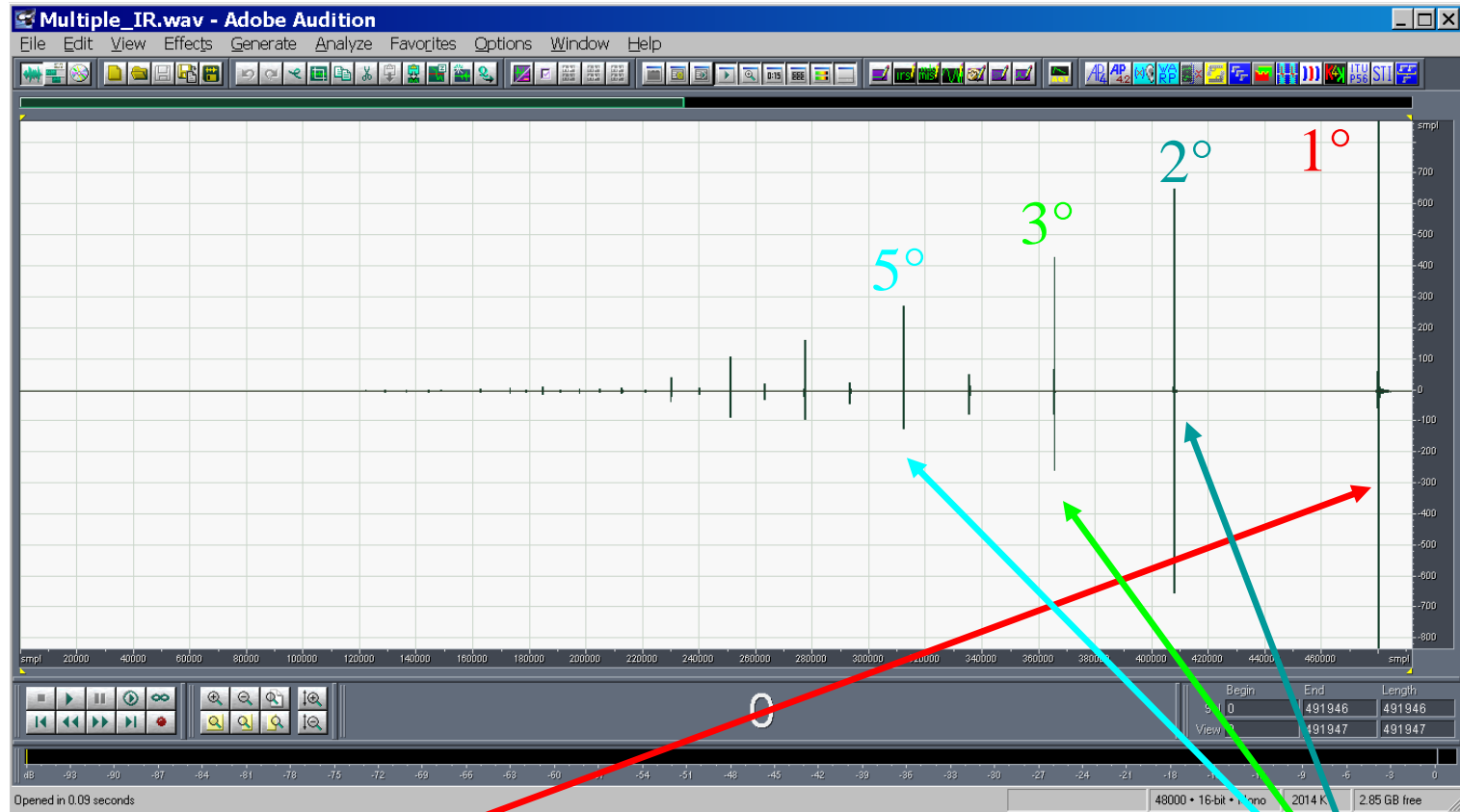
Test Signal  $x(t)$



Inverse Filter  $z(t)$



# Result of the deconvolution

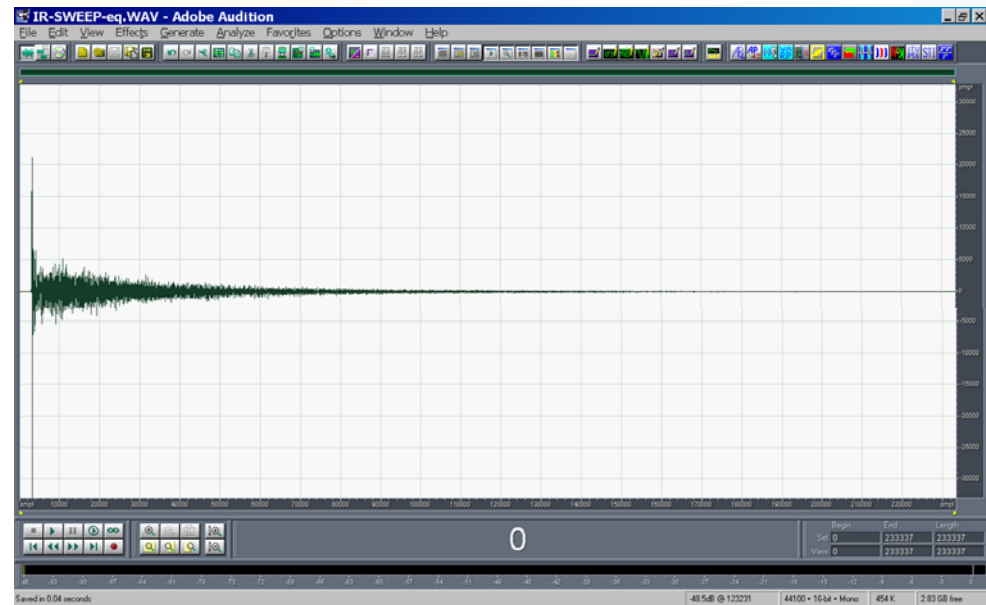
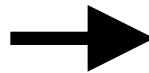
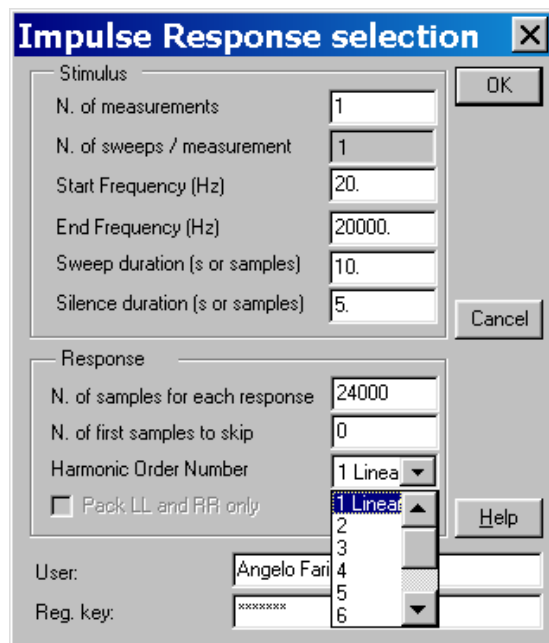


**The last** impulse response is the linear one, **the preceding** are the harmonics distortion products of various orders



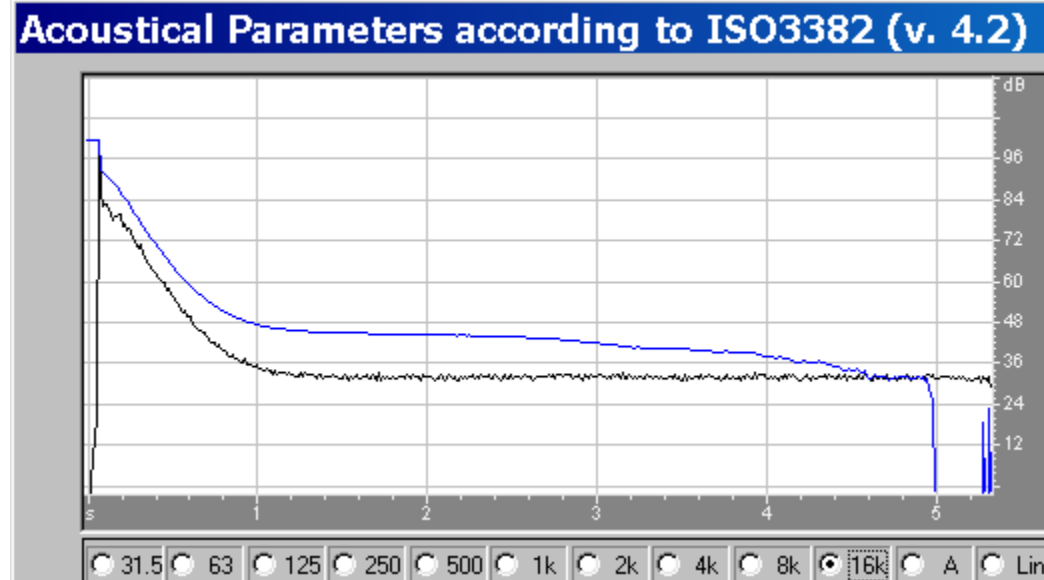
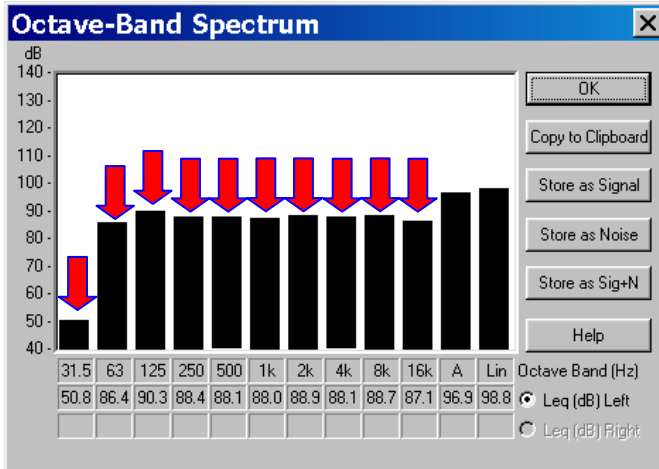
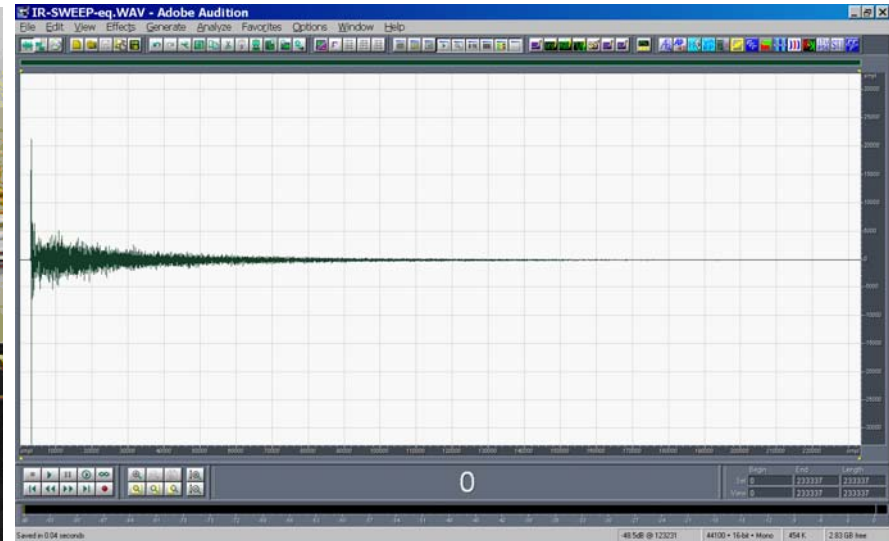
# IR Selection

- After the sequence of impulse responses has been obtained, it is possible to select and extract just one of them (the 1°-order - Linear in this example):



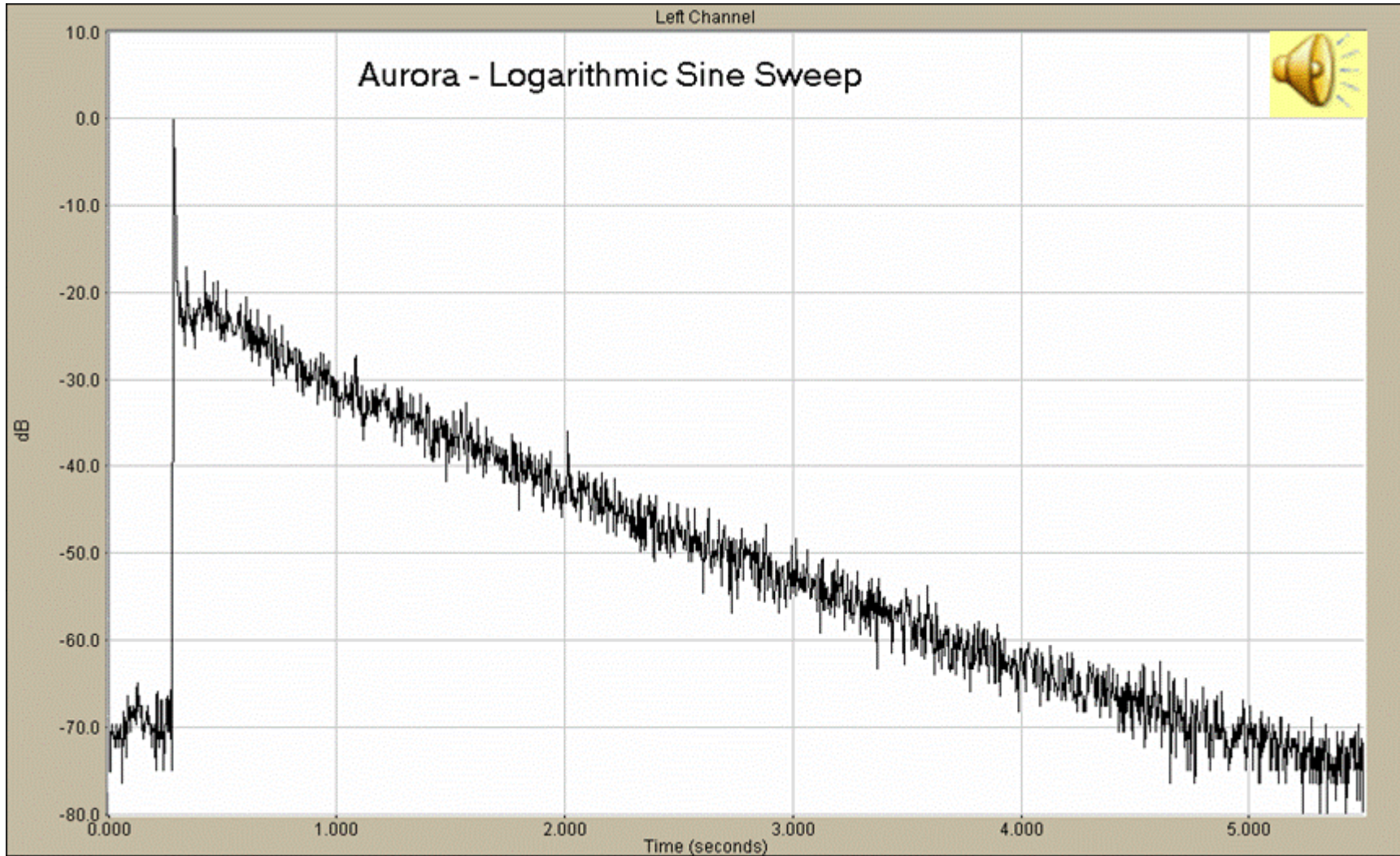


# Example of an ESS impulse response





# Maximum Length Sequence vs. Exp. Sine Sweep





# Post processing of impulse responses

- A special plugin has been developed for the computation of STI according to IEC-EN 60268-16:2003

### STI & Octave Band Analysis

Calibration (Octave Analysis)

Full Scale  Leq

Calibration value (dB):

Compute Octave Band Spectrum

Load SPL Values from File... Save SPL Values to File...

Hz	BackGnd Noise Level	Signal Level	Signal + Noise Level
125	48.0	70.9	70.9
250	45.0	70.9	70.9
500	42.0	67.2	67.2
1k	39.0	61.2	61.2
2k	36.0	55.2	55.3
4k	33.0	49.2	49.3
8k	30.0	43.2	43.4

Impulse Response Analysis

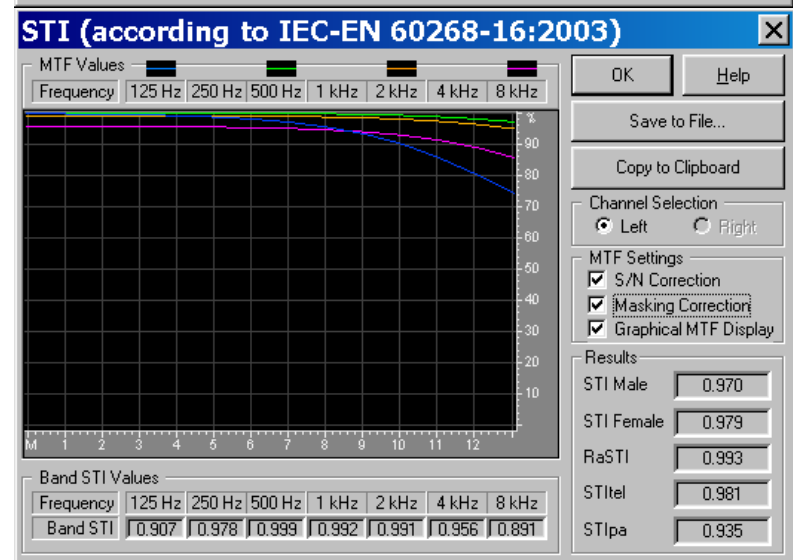
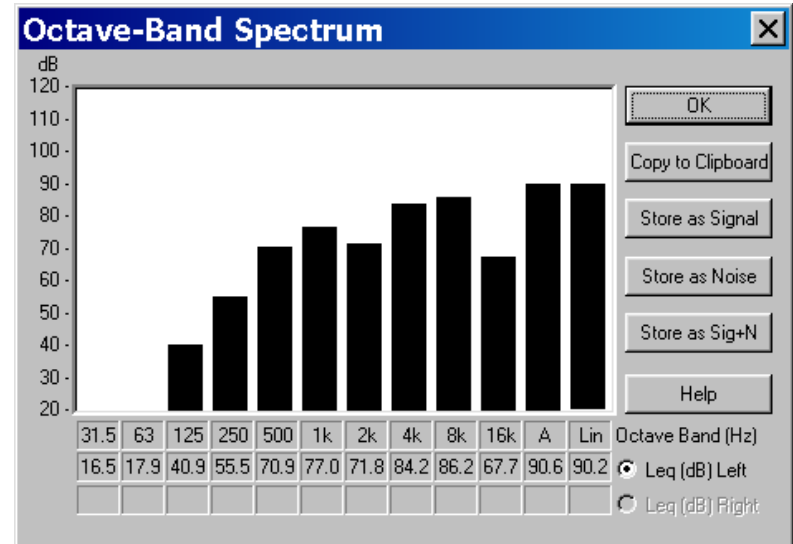
First Arrival Threshold (% of Full Scale):

Compute STI

User:

Reg. key:

Close Help







# Post processing of impulse responses

- A special plugin has been developed for performing analysis of acoustical parameters according to ISO-3382

**Acoustical Paramete...** [X]

User Defined Reverberation Time Extremes:  
( -5. dB , -15. dB )

Enable Noise Correction  
 EDT without linear regression

First Arrival Time Threshold (% of FS): 4  
Peak SPL value corresponding to FS: 120.

Stereo Mode

2 Omnidirectional Microphones  
 Soundfield Microphone (WY)  
 Omni/Eight microphone  
 p-p Sound Intensity Probe

d (mm): 12.0 c (m/s): 340.0

Binaural Dummy Head

IACC Integration [v]

User: Angelo Farina  
Reg. key: \*\*\*\*\*

[OK] [Cancel] [Help]

**Acoustical Parameters according to ISO3382 (v. 4.2)** [X]

[Close] [Help]

OK - keep processed  
Save Results to File...  
Copy Results to Clipboard  
Store G Reference Signal

Channel:  
 Left  Right

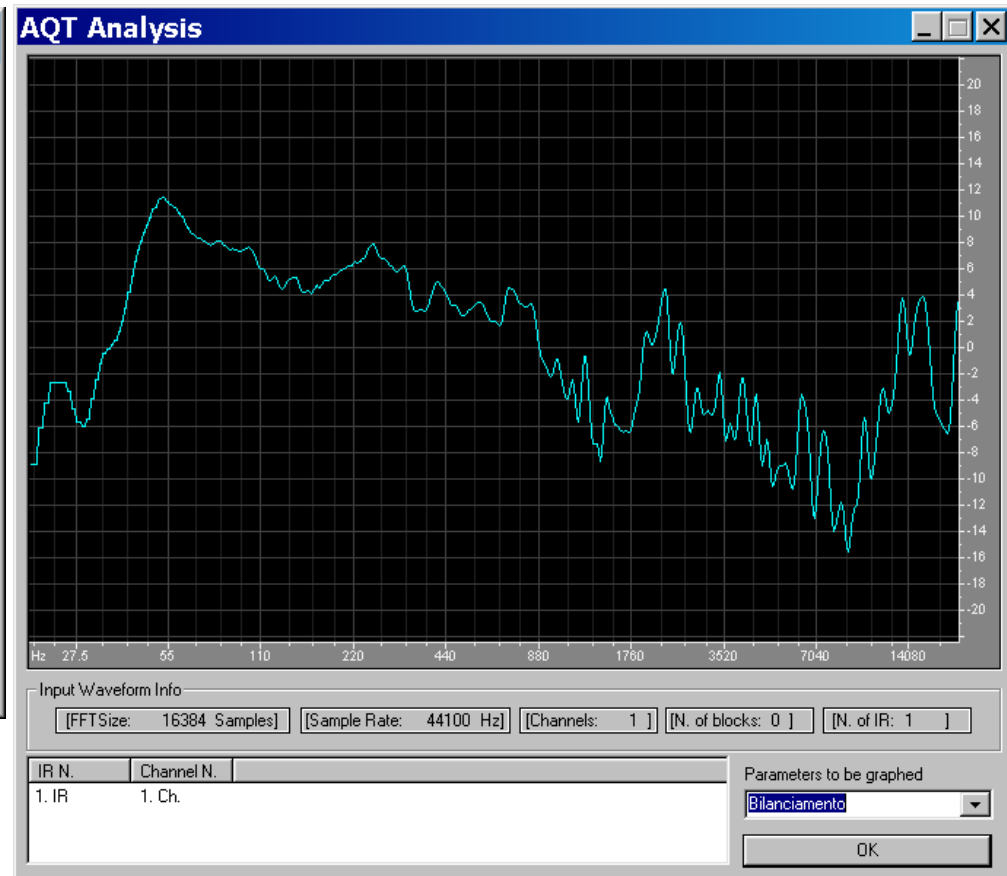
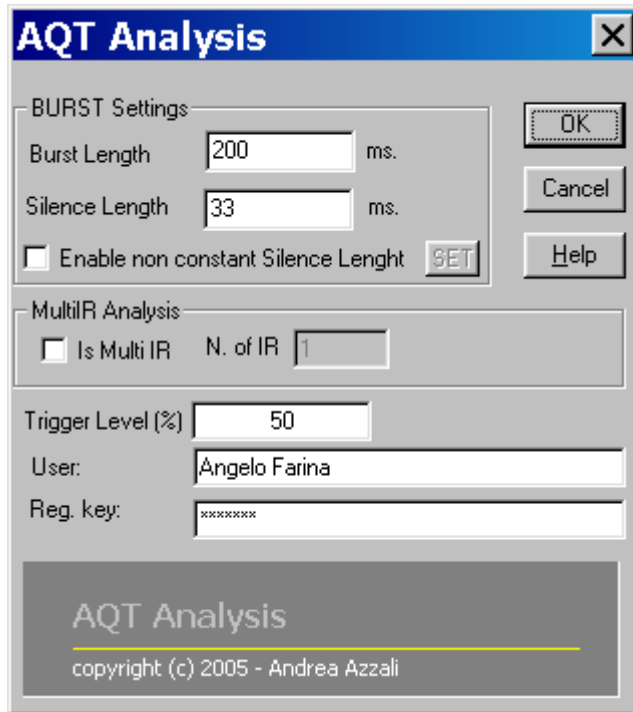
Tuser limits:  
(-5. dB, -15. dB)

31.5	63	125	250	500	1k	2k	4k	8k	16k	A	Lin	Freq. (Hz)
23.70	58.39	61.53	65.32	66.44	72.30	80.55	78.61	79.40	80.73	85.87	85.62	Signal (dB)
24.44	31.94	24.31	21.28	18.96	20.63	24.14	26.10	32.18	37.87	35.85	38.67	Noise (dB)
-45.30	-10.61	-7.47	-3.68	-2.56	3.30	11.55	9.61	10.40	11.73	8.87	8.62	G (dB)
-2.47	-3.41	-2.95	-6.29	-4.08	-5.01	-4.08	-1.32	5.90	9.34	-0.80	0.11	C50 (dB)
-0.63	-1.23	-1.21	-4.58	-2.74	-2.71	-1.69	0.91	8.42	12.48	1.12	1.92	C80 (dB)
36.17	31.33	33.66	19.04	28.12	23.96	28.10	42.48	79.57	89.58	45.41	50.64	D50 (%)
204.39	163.99	196.26	189.81	170.83	161.80	150.38	113.13	32.48	22.26	110.19	99.58	Ts (ms)
4.48	1.93	2.82	2.24	2.21	2.16	2.03	1.73	0.68	0.31	1.82	1.76	EDT (s)
--	3.00	3.07	2.06	2.14	2.26	2.14	1.82	0.84	0.56	2.01	1.99	Tuser (s)
--	2.76	2.84	2.32	2.15	2.24	2.16	1.92	0.99	0.60	2.07	2.07	T20 (s)
--	2.87	2.77	2.51	2.14	2.27	2.20	2.00	1.04	0.65	2.13	2.15	T30 (s)
1.00	1.00	1.00	0.97	0.51	0.40	0.42	0.55	0.58	0.57	0.50	0.52	IACC (Early)
-0.02	-0.02	-0.05	-0.02	-0.09	-0.07	-0.05	-0.02	-0.05	-0.02	-0.02	-0.02	t IACC (ms)
1.86	1.79	1.11	0.57	0.27	0.16	0.09	0.07	0.07	0.05	0.07	0.07	w IACC (ms)



# The new AQT plugin for Audition

- The new module is still under development and will allow for very fast computation of the AQT (Dynamic Frequency Response) curve from within Adobe Audition





# Spatial analysis by directive microphones

- The initial approach was to use directive microphones for gathering some information about the spatial properties of the sound field “as perceived by the listener”
- Two apparently different approaches emerged: binaural dummy heads and pressure-velocity microphones:



Binaural  
microphone (left)

and

variable-directivity  
microphone (right)



# IACC “objective” spatial parameter

- It was attempted to “quantify” the “spatiality” of a room by means of “objective” parameters, based on 2-channels impulse responses measured with directive microphones
- The most famous “spatial” parameter is IACC (Inter Aural Cross Correlation), based on binaural IR measurements

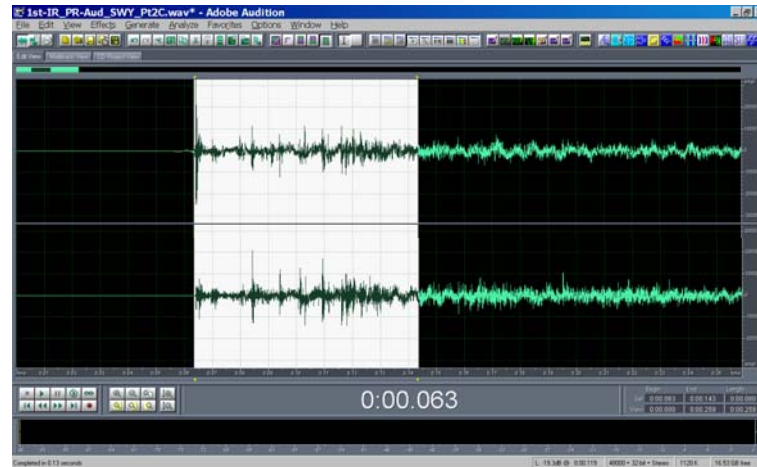
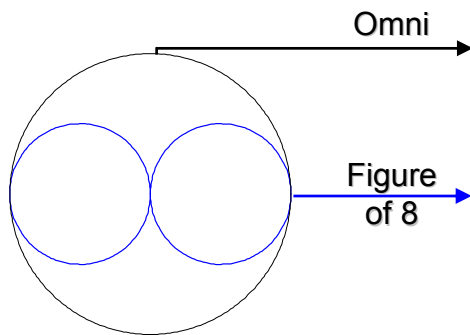


$$\rho(t) = \frac{\int_0^{80\text{ms}} p_L(\tau) \cdot p_R(\tau+t) \cdot d\tau}{\sqrt{\int_0^{80\text{ms}} p_L^2(\tau) \cdot d\tau \cdot \int_0^{80\text{ms}} p_R^2(\tau+t) \cdot d\tau}}$$

$$IACC_E = \text{Max}[\rho(t)] \quad t \in [-1\text{ms} \dots +1\text{ms}]$$

# LF “objective” spatial parameter

- Other “spatial” parameters are the Lateral Energy ratio LF
- This is defined from a 2-channels impulse response, the first channel is a standard omni microphone, the second channel is a “figure-of-eight” microphone:



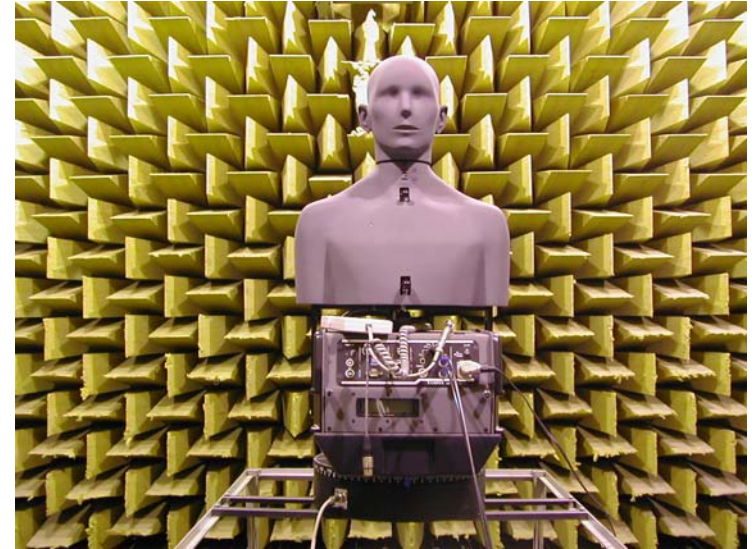
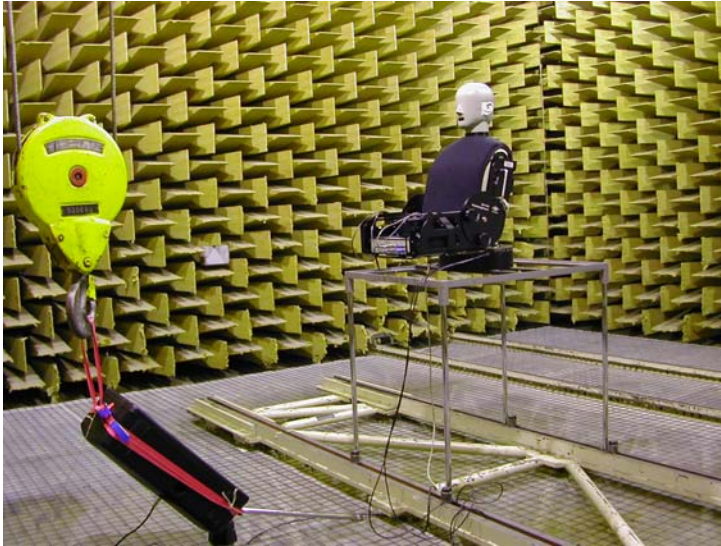
$$LF = \frac{\int_{0ms}^{80ms} h_8^2(\tau) \cdot d\tau}{\int_{0ms}^{80ms} h_o^2(\tau) \cdot d\tau}$$





# Are IACC measurements reproducible?

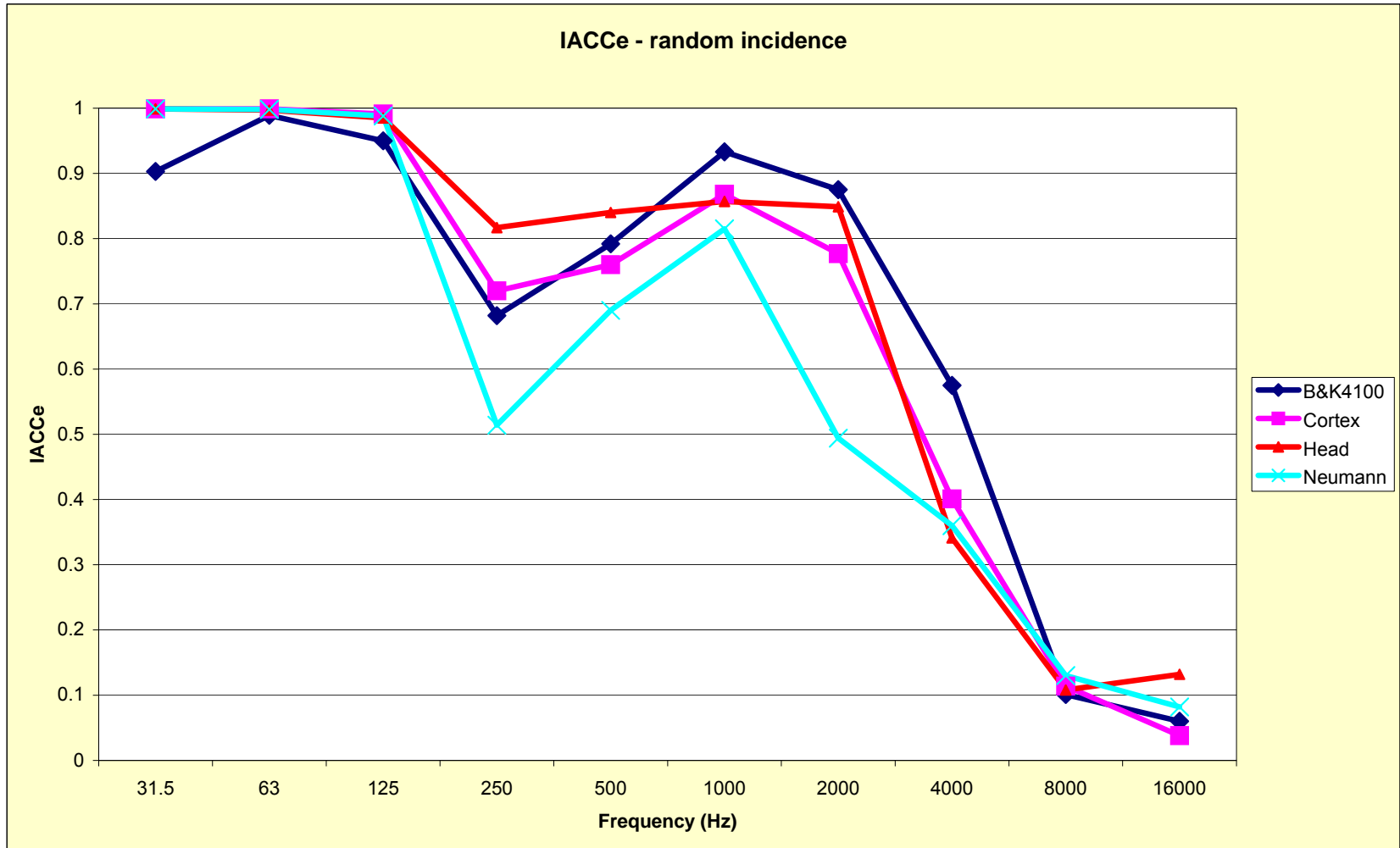
- Experiment performed in anechoic room - same loudspeaker, same source and receiver positions, 5 binaural dummy heads





# Are IACC measurements reproducible?

- Diffuse field - huge difference among the 4 dummy heads





# Are LF measurements reproducible?

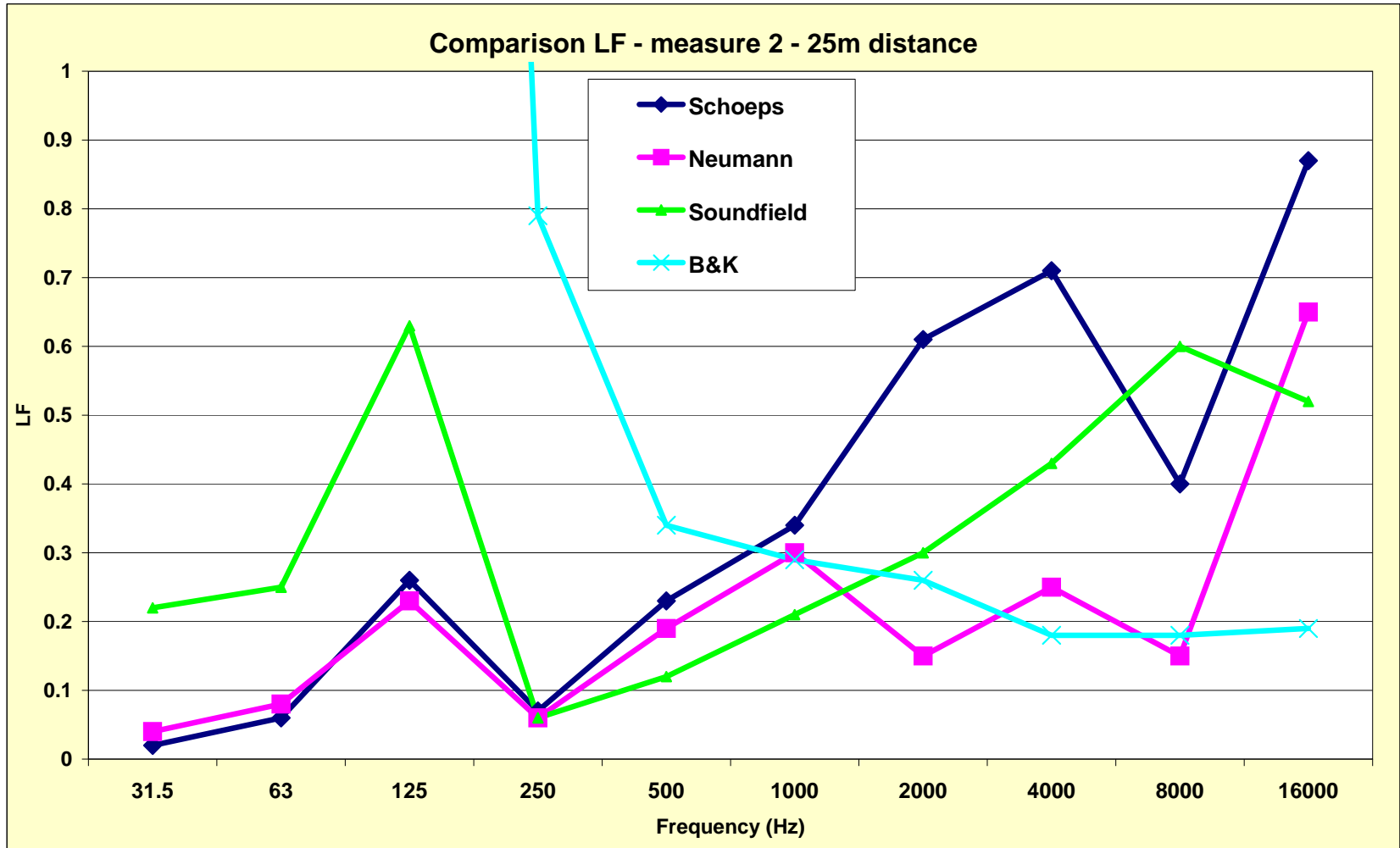
- Experiment performed in the Auditorium of Parma - same loudspeaker, same source and receiver positions, 4 pressure-velocity microphones





# Are LF measurements reproducible?

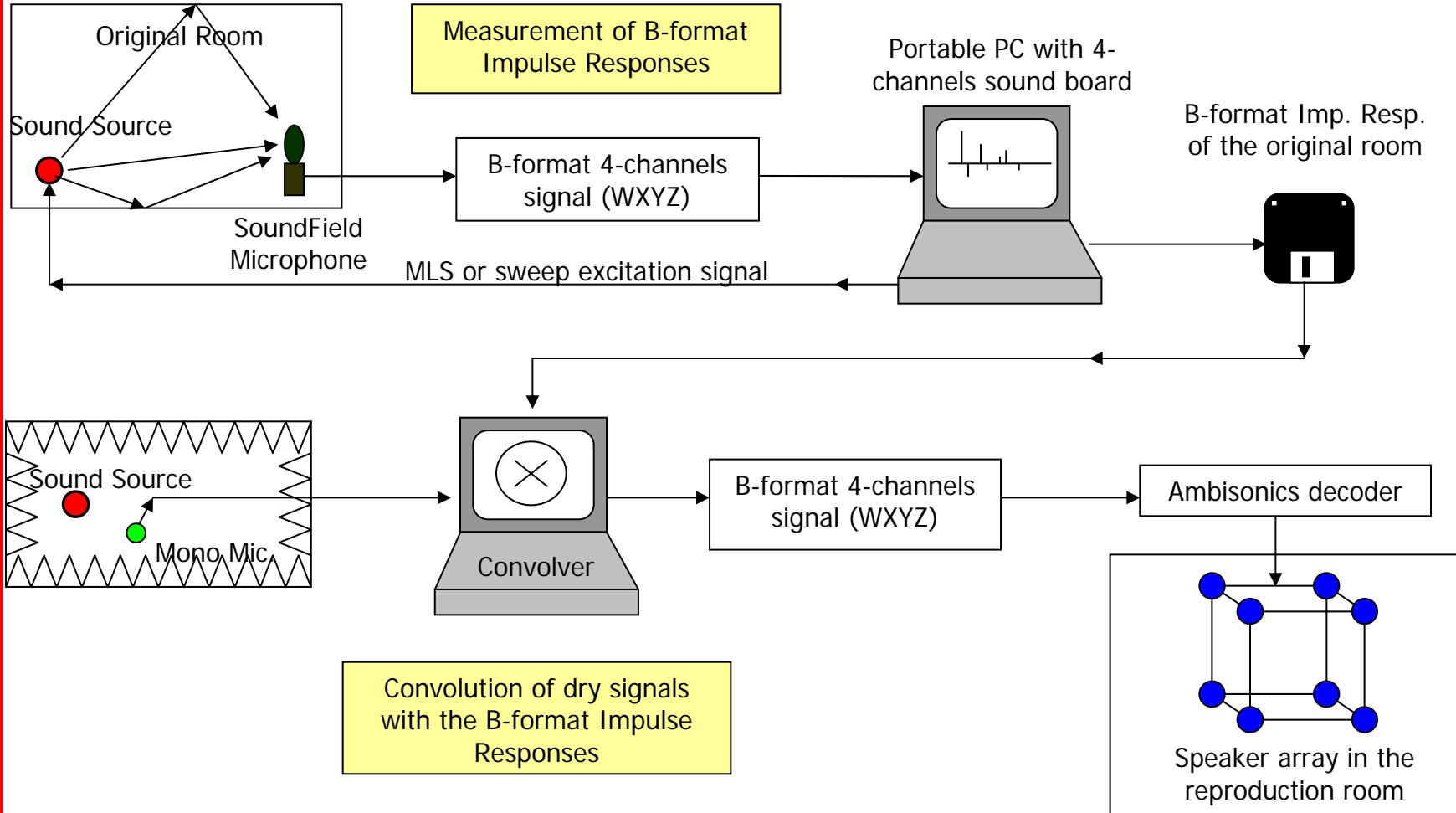
- At 25 m distance, the scatter is really big







# 3D Impulse Response (Gerzon, 1975)

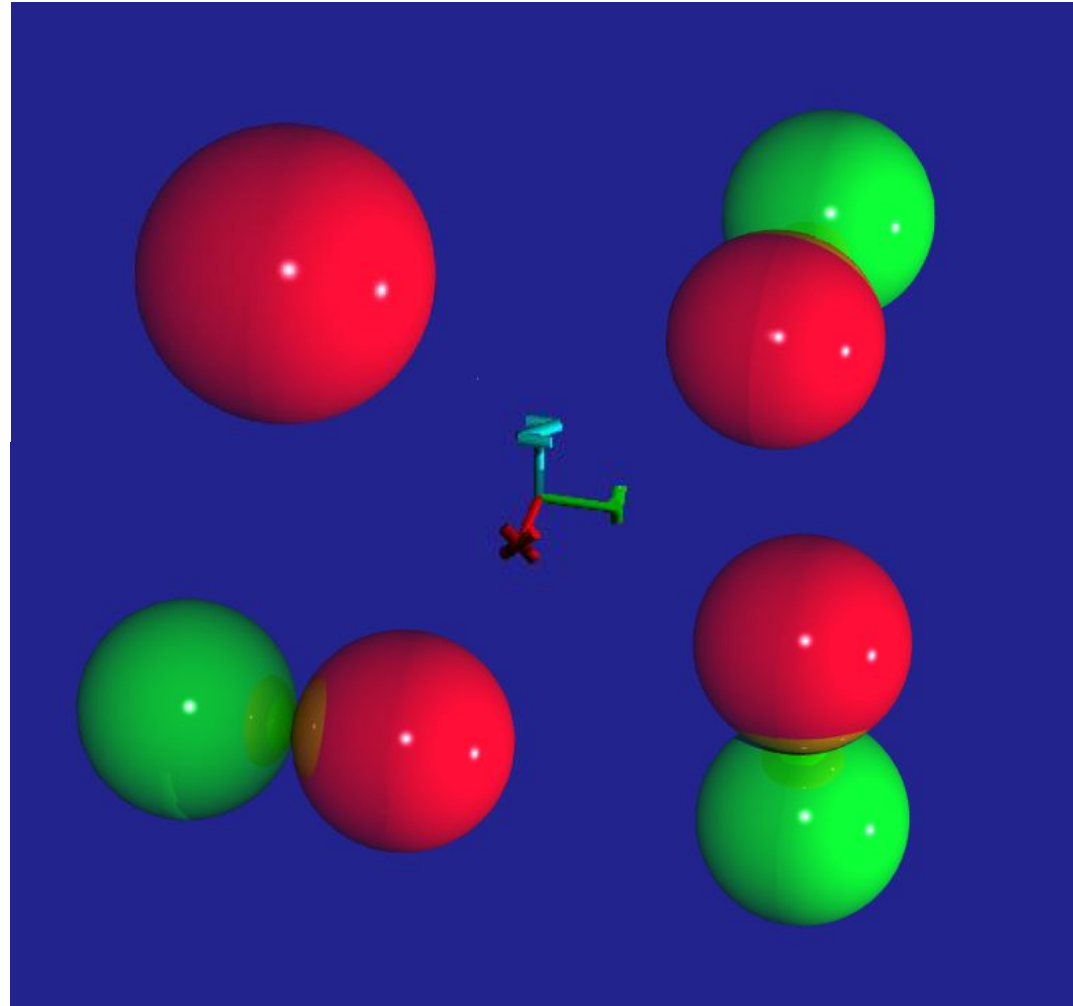






# 3D extension of the pressure-velocity measurements

- The Soundfield microphone allows for simultaneous measurements of the omnidirectional pressure and of the three cartesian components of particle velocity (figure-of-8 patterns)





# The Waves project (2003)

- The original idea of Michael Gerzon was finally put in practice in 2003, thanks to the Israeli-based company WAVES
- More than 50 theatres all around the world were measured, capturing 3D IRs (4-channels B-format with a Soundfield microphone)
- The measurements did also include binaural impulse responses, and a circular-array of microphone positions
- More details on [WWW.ACOUSTICS.NET](http://WWW.ACOUSTICS.NET)





# The Future



# The Future

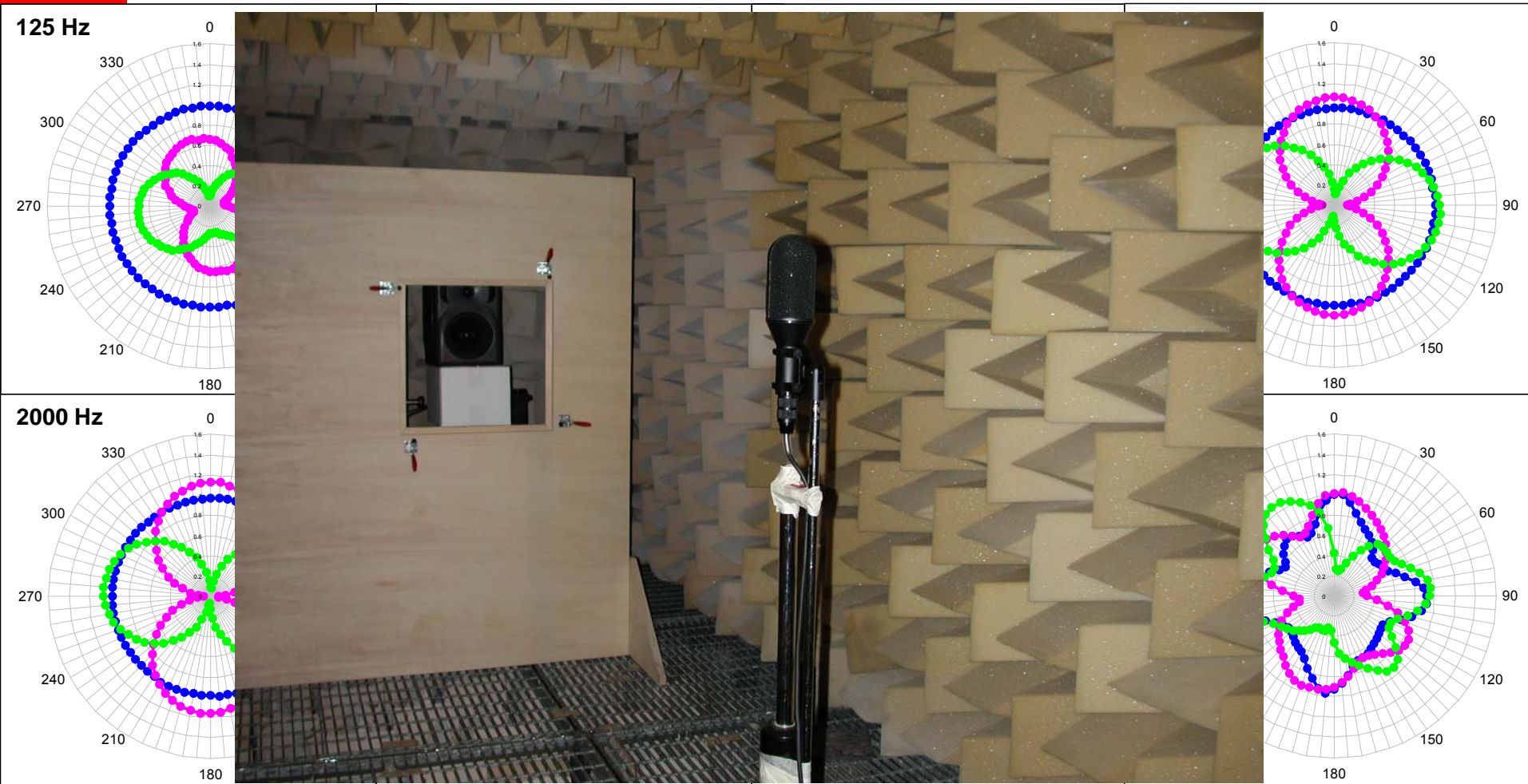
- **Microphone arrays capable of synthesizing arbitrary directivity patterns**
- **Advanced spatial analysis of the sound field employing spherical harmonics (Ambisonics - 1<sup>o</sup> order or higher)**
- **Loudspeaker arrays capable of synthesizing arbitrary directivity patterns**
- **Generalized solution in which both the directivities of the source and of the receiver are represented as a spherical harmonics expansion**



# Directivity of transducers



## Soundfield ST-250 microphone

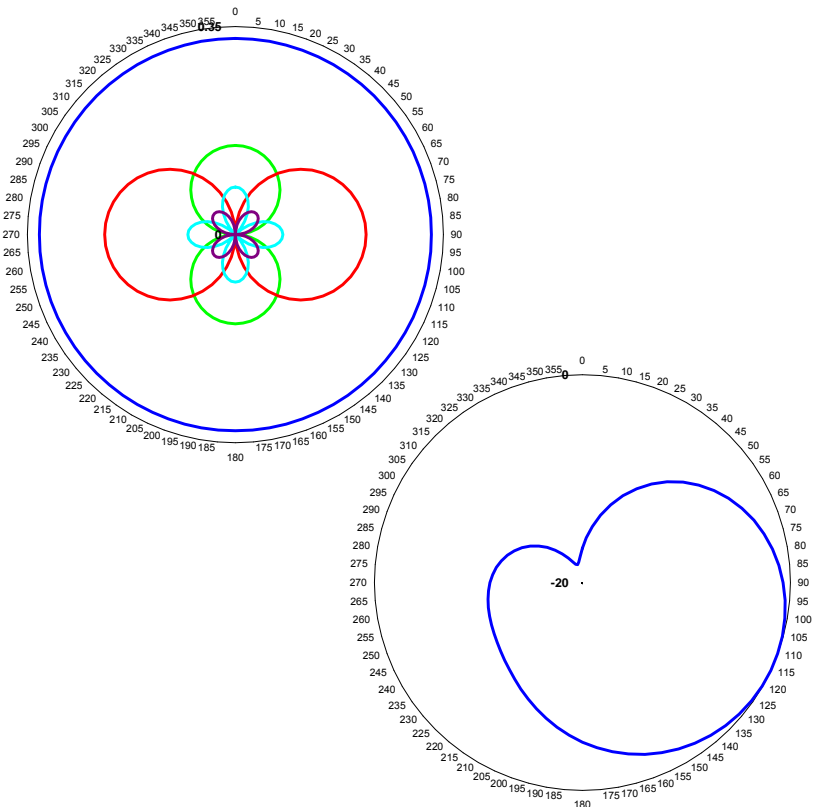
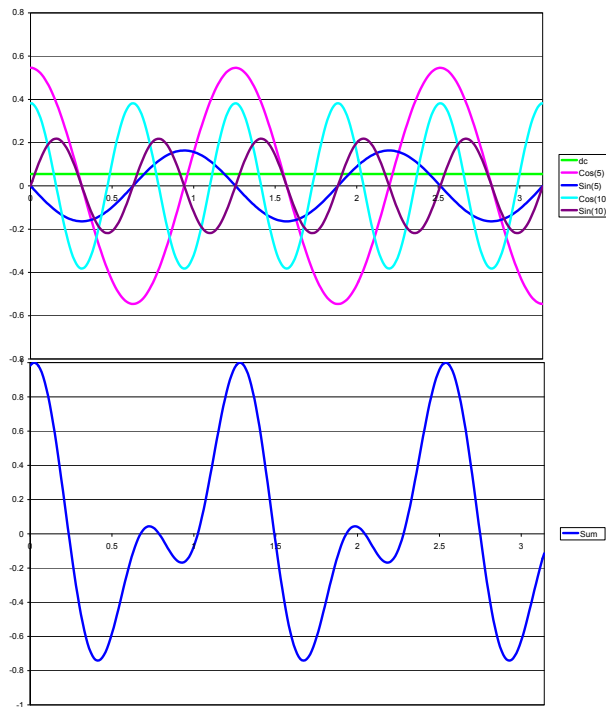






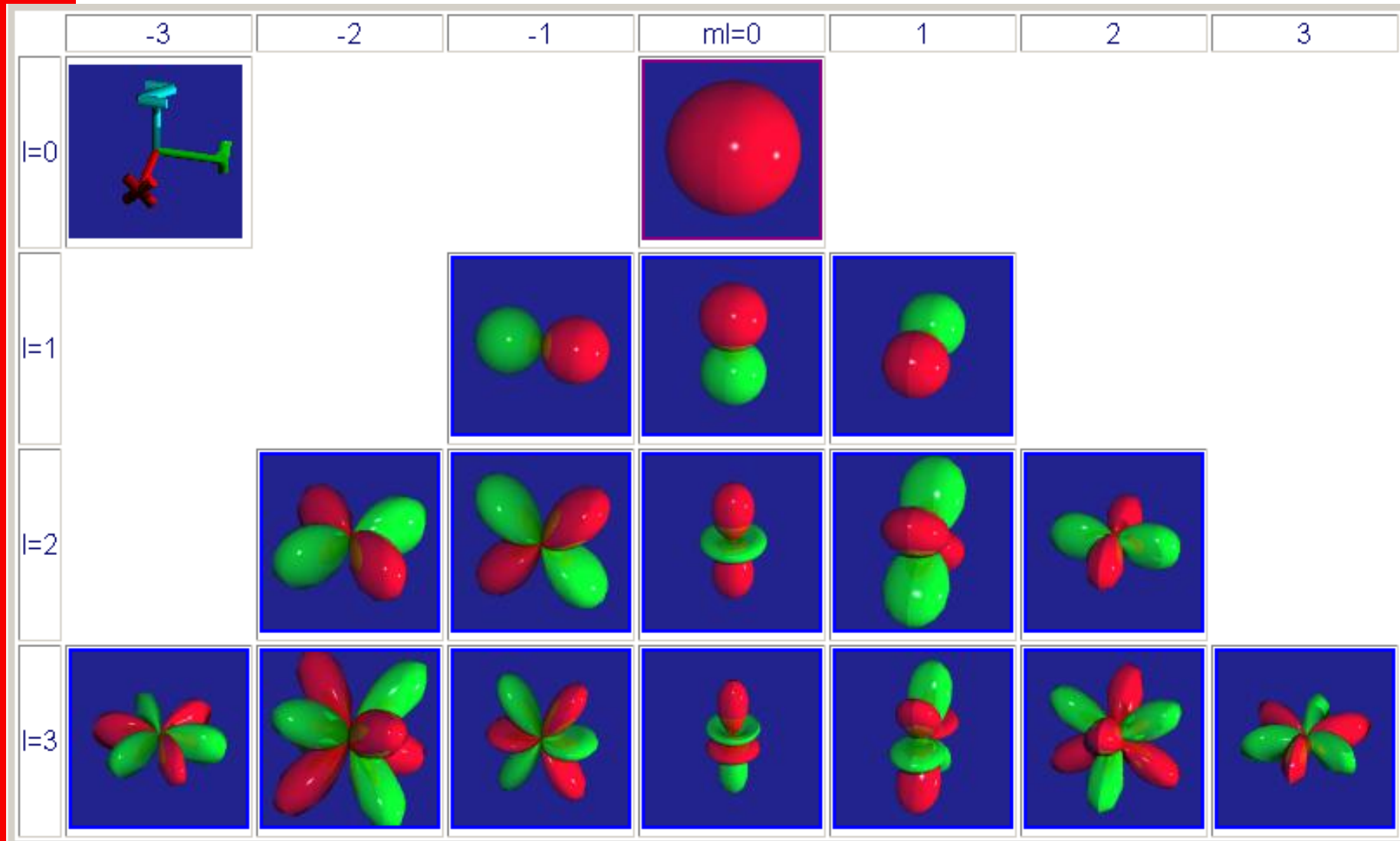
# How to get better spatial resolution?

- The answer is simple: analyze the spatial distribution of both source and receiver by means of higher-order spherical harmonics expansion
- Spherical harmonics analysis is the equivalent, in space domain, of the Fourier analysis in time domain
- As a complex time-domain waveform can be thought as the sum of a number of sinusoidal and cosinusoidal functions, so a complex spatial distribution around a given notional point can be expressed as the sum of a number of spherical harmonic functions





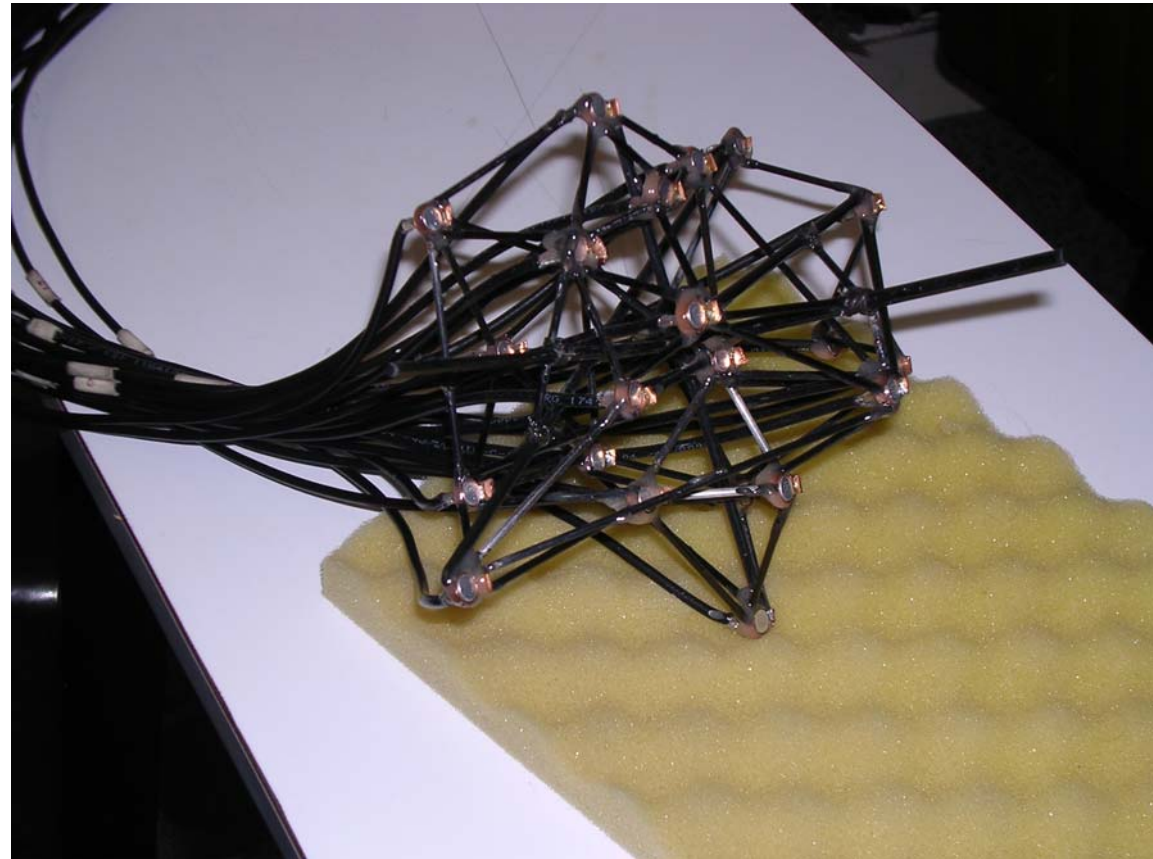
# Higher-order spherical harmonics expansion



# 3°-order microphone (Trinnov - France)



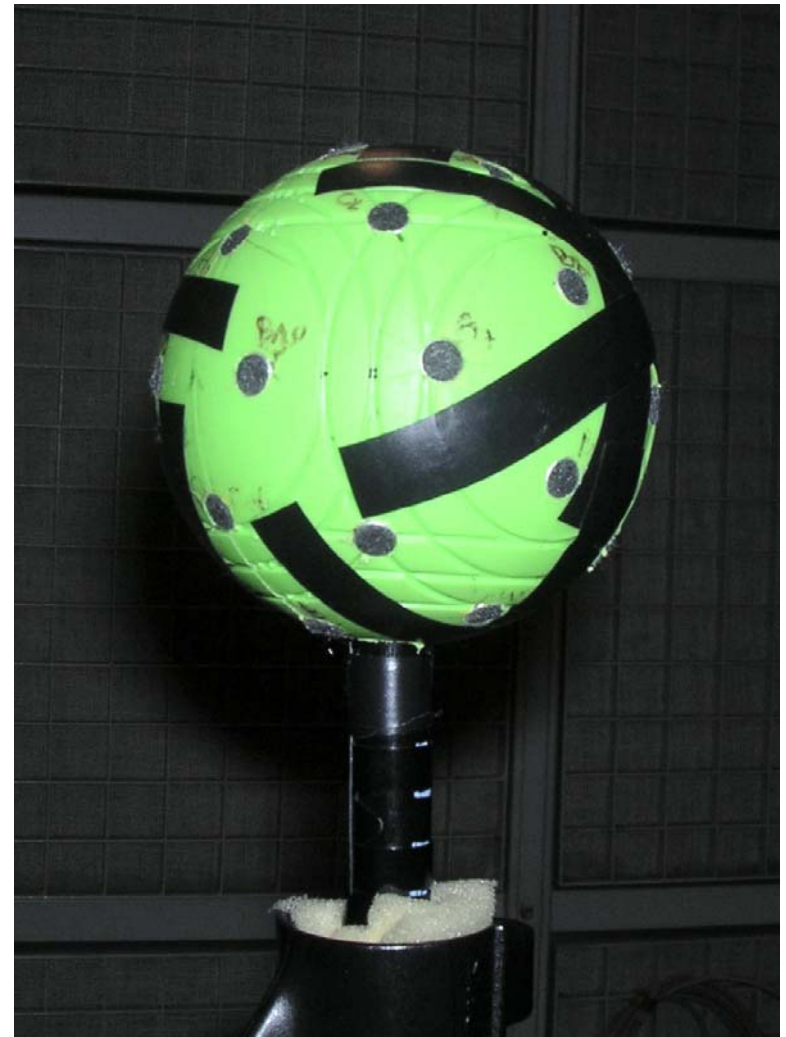
- Arnoud Laborie developed a 24-capsule compact microphone array - by means of advanced digital filtering, spherical ahrmonic signals up to 3° order are obtained (16 channels)



# 4°-order microphone (France Telecom)



- Jerome Daniel and Sebastien Moreau built samples of 32-capsules spherical arrays - these allow for extractions of microphone signals up to 4° order (25 channels)





# 4°-order microphone (Italy)



- Angelo Farina's spherical mike (32 capsules)





# 5°-order microphones (University of Sydney)



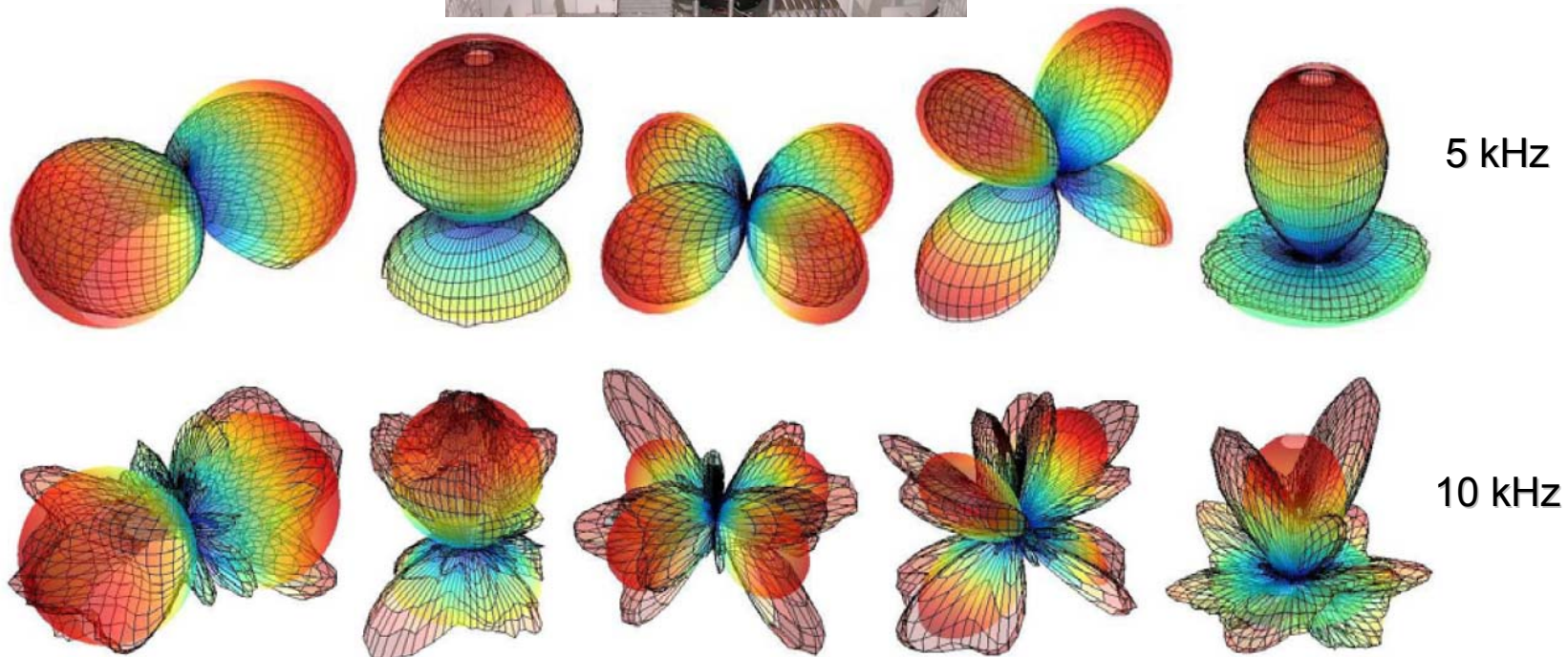
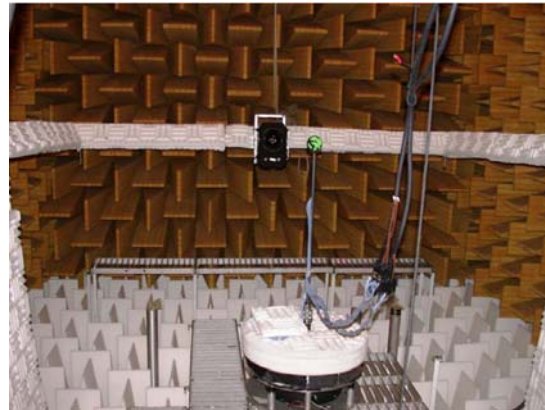
- Chris Craig's dual-sphere concentric mike (64 capsules)
- And his 32-capsules cylindrical mike



# Verification of high-order patterns



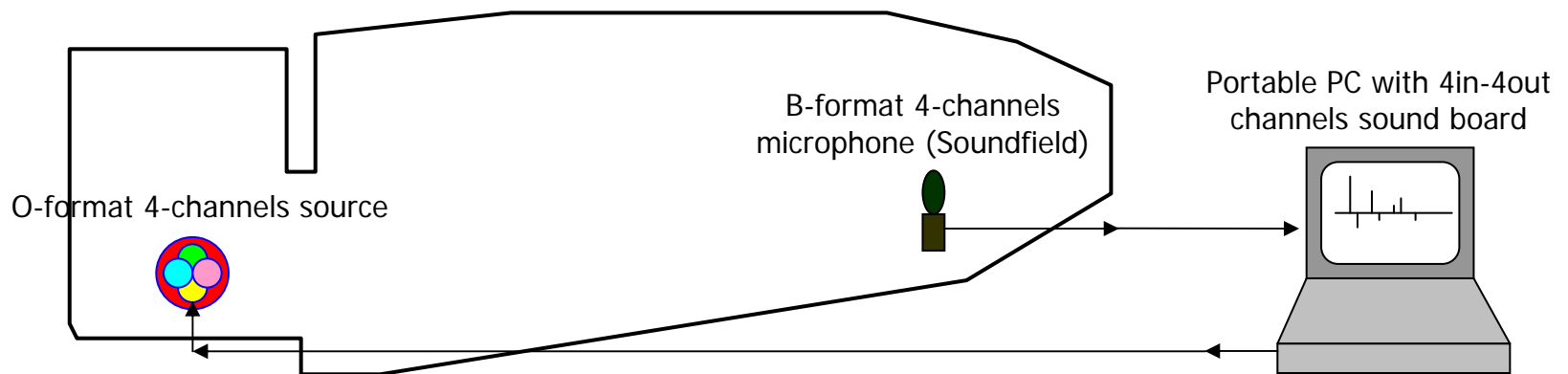
- Sebastien Moreau and Olivier Warusfel verified the directivity patterns of their 4<sup>o</sup>-order microphone array in the anechoic room of IRCAM (Paris)





# What about source directivity ?

- Current 3D IR sampling is still based on the usage of an “omnidirectional” source
- The knowledge of the 3D IR measured in this way provide no information about the soundfield generated inside the room from a directive source (i.e., a musical instrument, a singer, etc.)
- Dave Malham suggested to represent also the source directivity with a set of spherical harmonics, called O-format - this is perfectly reciprocal to the representation of the microphone directivity with the B-format signals (Soundfield microphone).
- Consequently, a complete and reciprocal spatial transfer function can be defined, employing a 4-channels O-format source and a 4-channels B-format receiver:

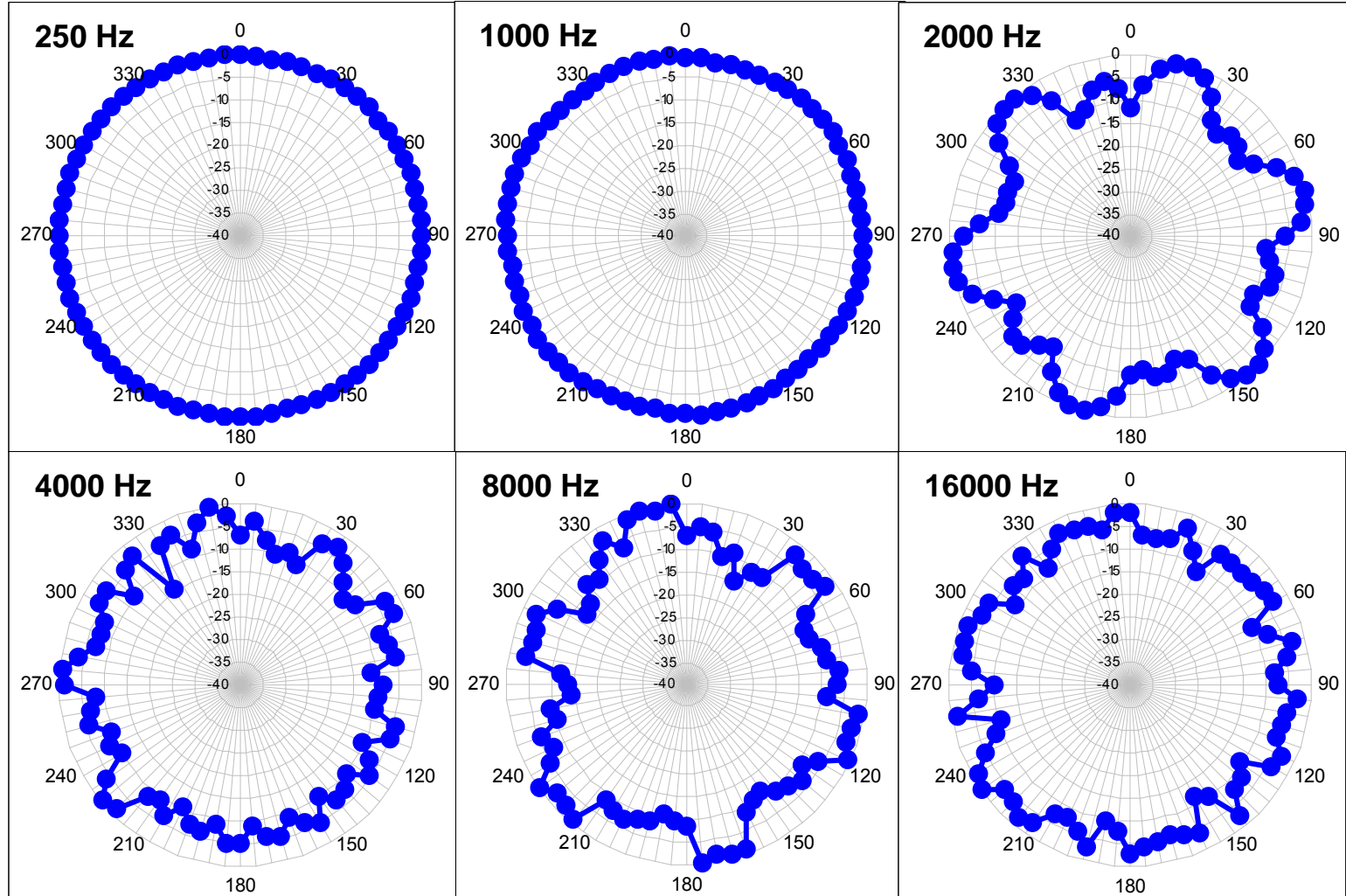




# Directivity of transducers



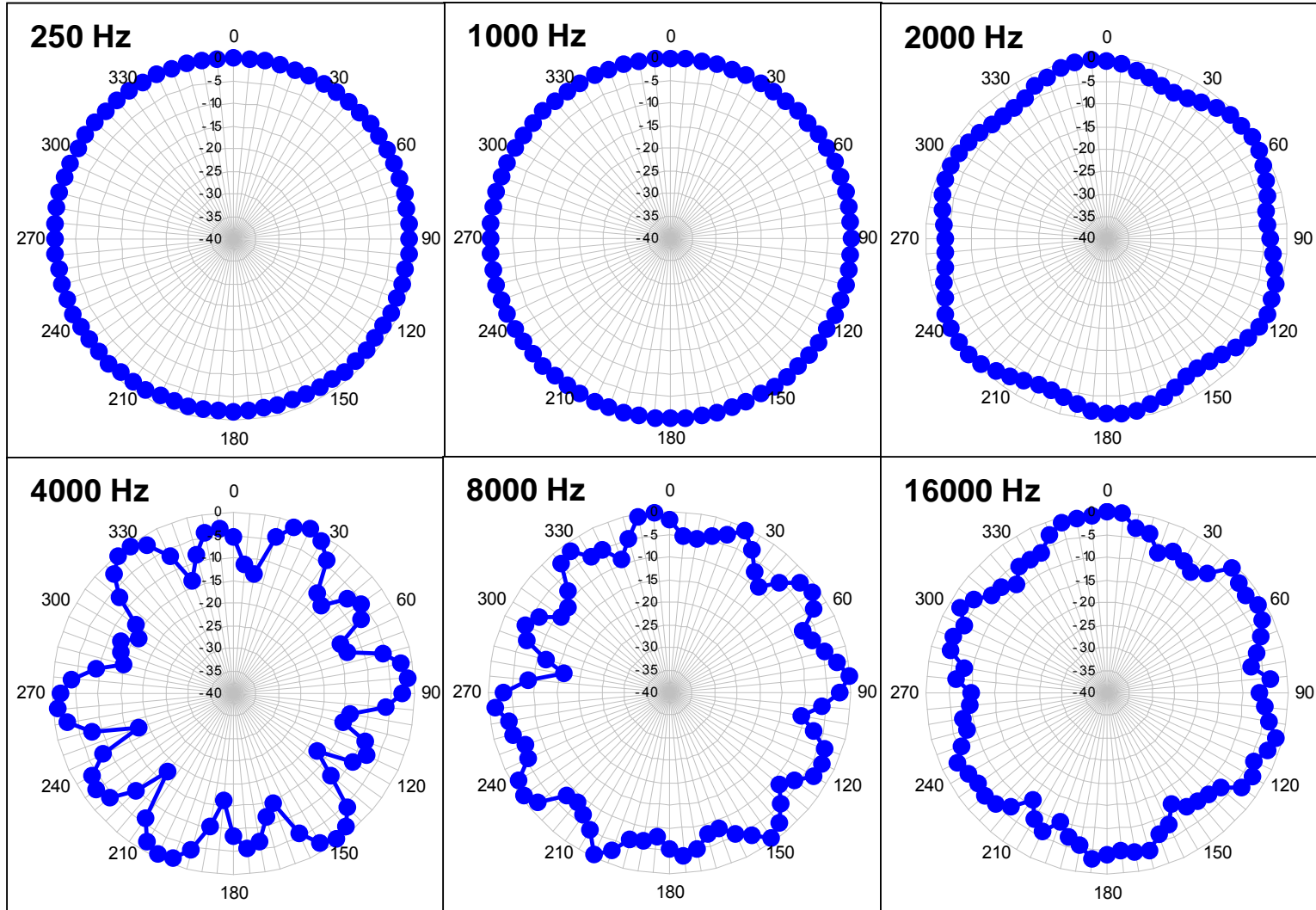
## LookLine D-300 dodechaedron



# Directivity of transducers



## LookLine D-200 dodechaedron

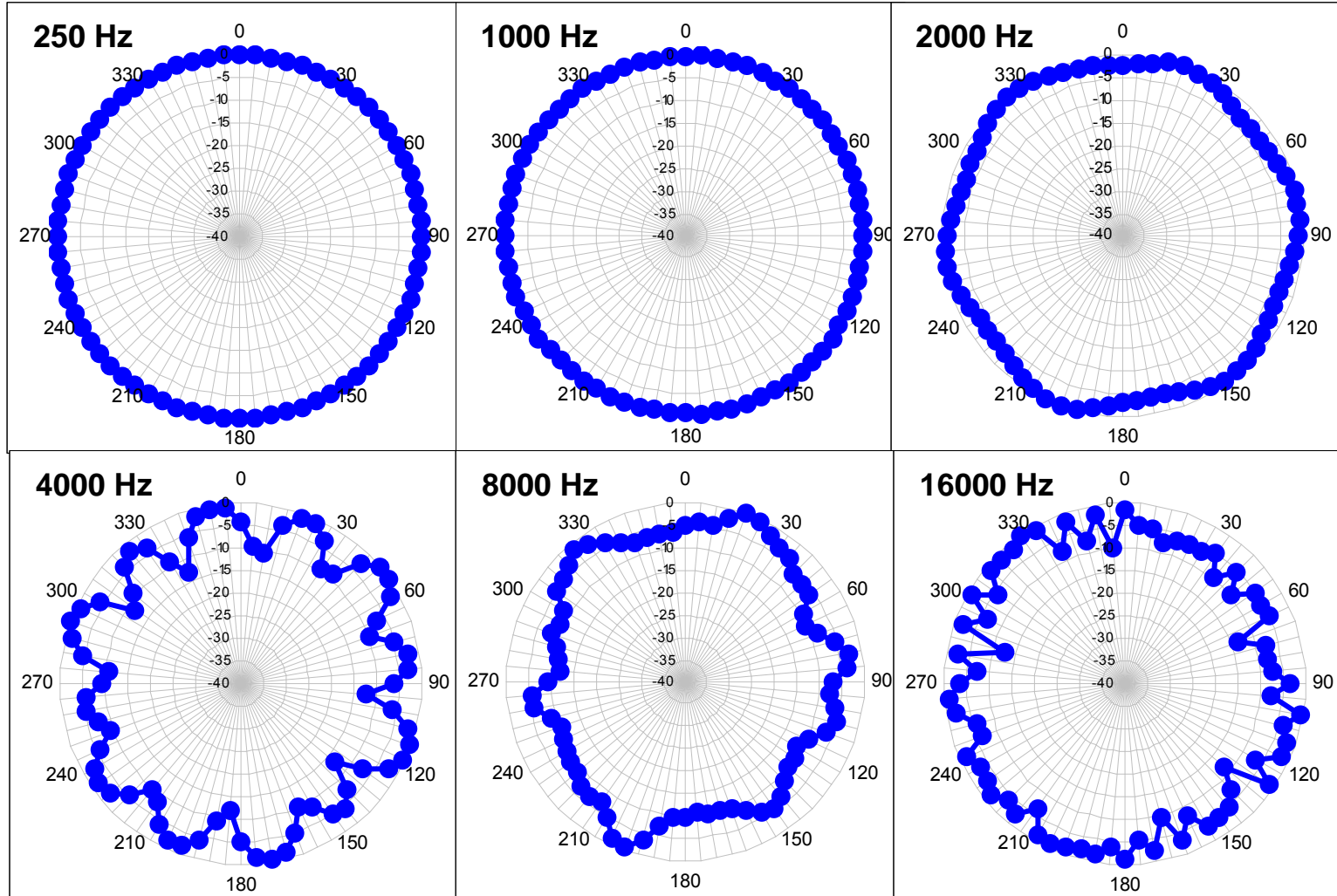




# Directivity of transducers



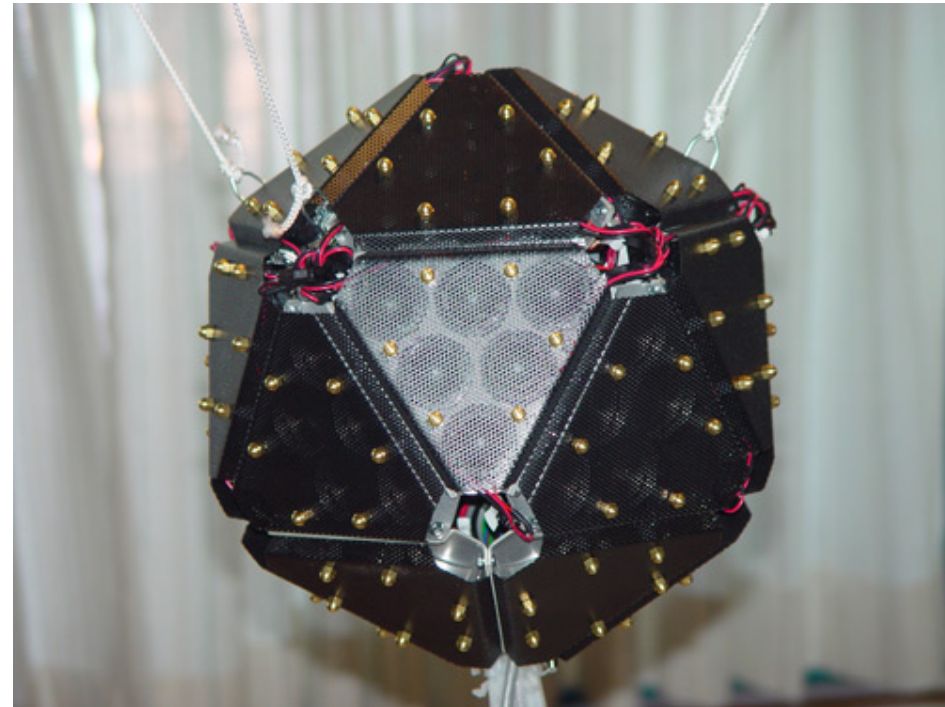
## Omnisonic 1000 dodechaedron





# High-order sound source

- Adrian Freed, Peter Kassakian, and David Wessel (CNMAT) developed a new 120-loudspeakers, digitally controlled sound source, capable of synthesizing sound emission according to spherical harmonics patterns up to 5<sup>o</sup> order.



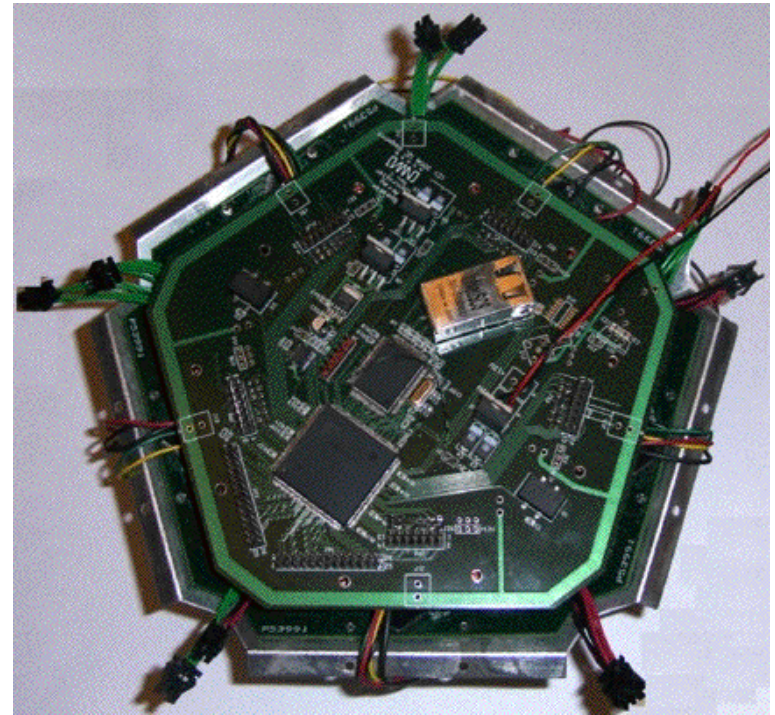




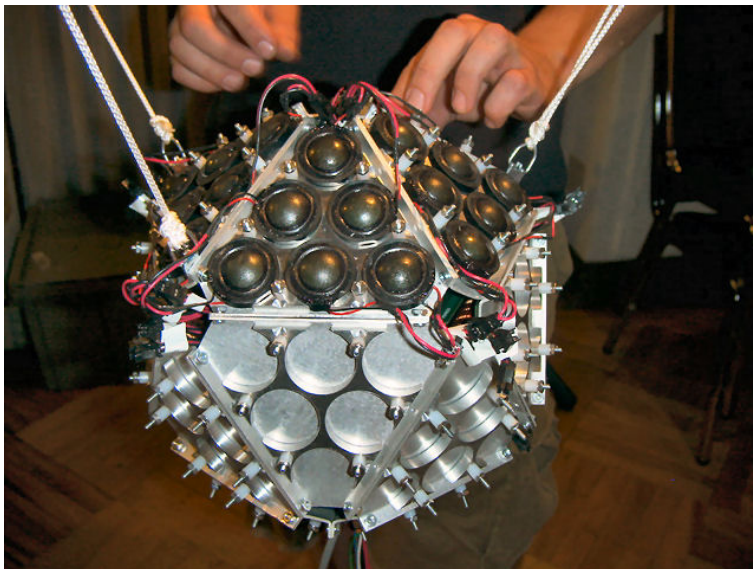
# Technical details of high-order source



- **Class-D embedded amplifiers**



- **Embedded ethernet interface and DSP processing**

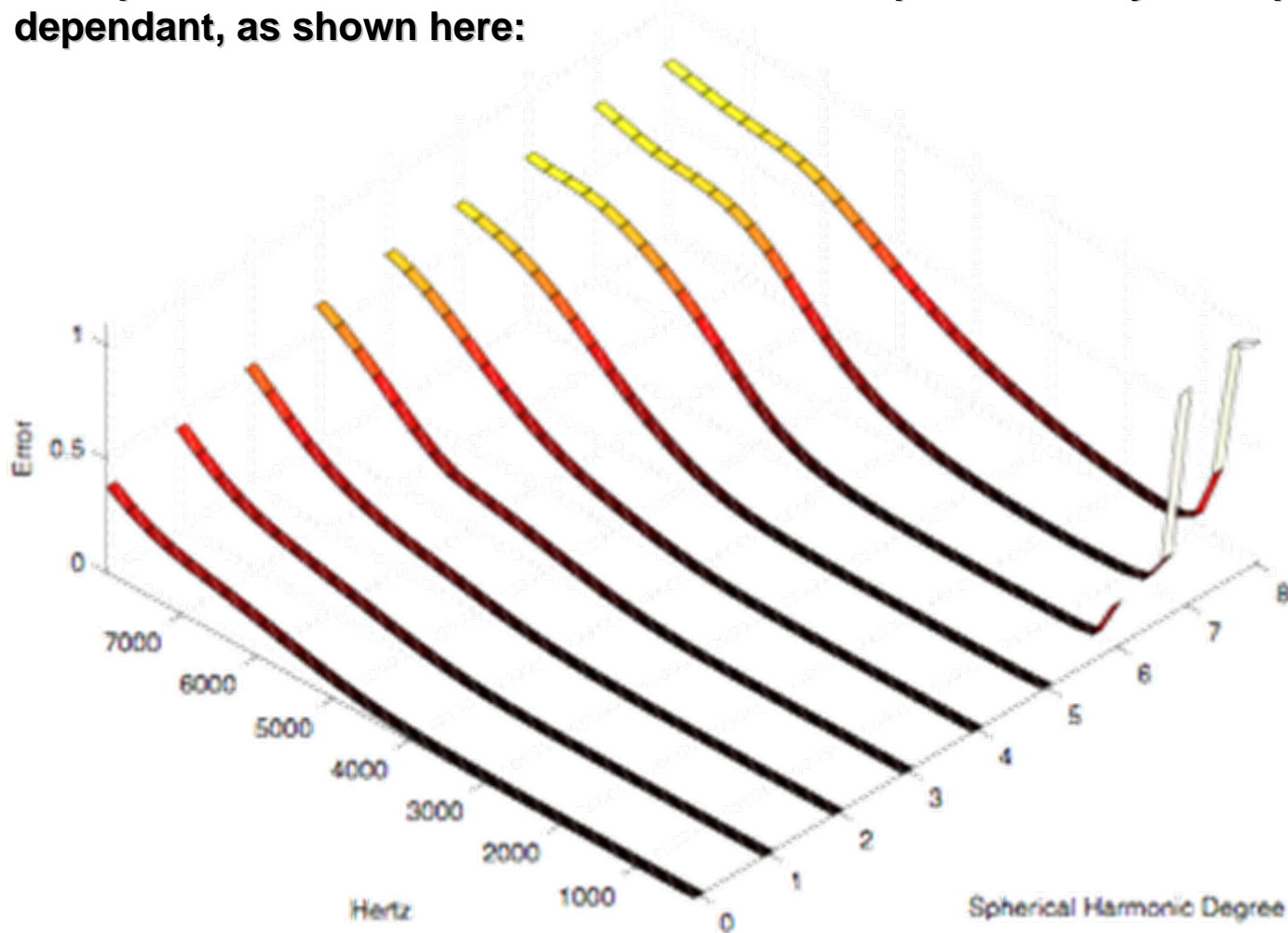


- **Long-excursion special Meyer Sound drivers**



# Accuracy of spatial synthesis

- The spatial reconstruction error of a 120-loudspeakers array is frequency dependant, as shown here:

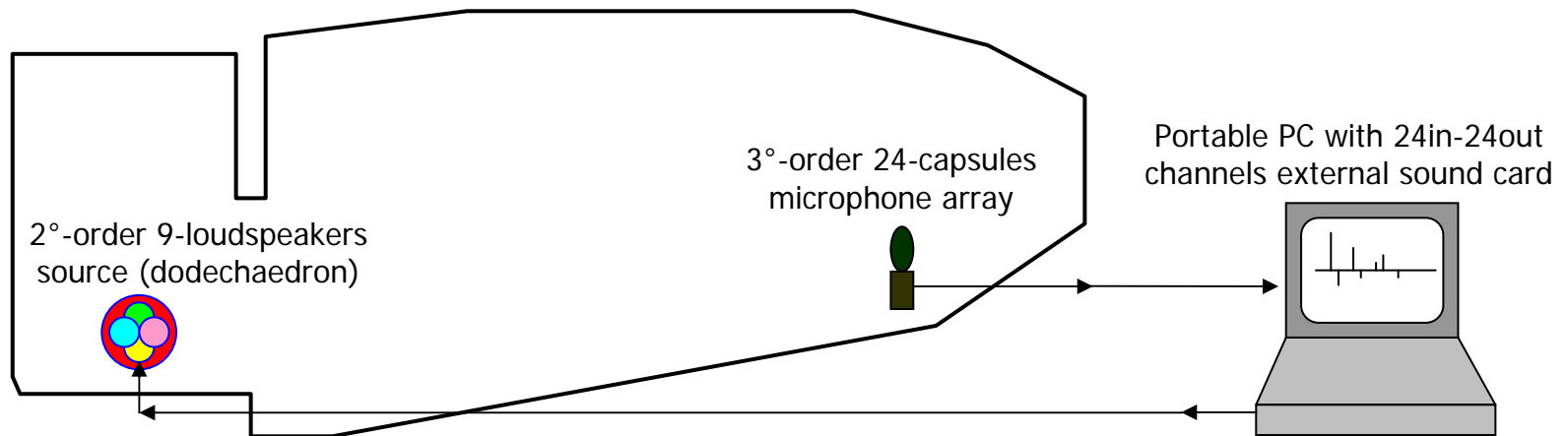


- The error is acceptably low over an extended frequency range up to 5<sup>o</sup>-order



# Complete high-order MIMO method

- Employing massive arrays of transducers, it will be feasible to sample the acoustical temporal-spatial transfer function of a room
- Currently available hardware and software tools make this practical only up to 4° order, which means 25 inputs and 25 outputs
- A complete measurement for a given source-receiver position pair takes approximately 10 minutes (25 sine sweeps of 15s each are generated one after the other, while all the microphone signals are sampled simultaneously)
- However, it has been seen that real-world sources can be already approximated quite well with 2°-order functions, and even the human HRTF directivities are reasonably approximated with 3°-order functions.







# Improving the ESS method



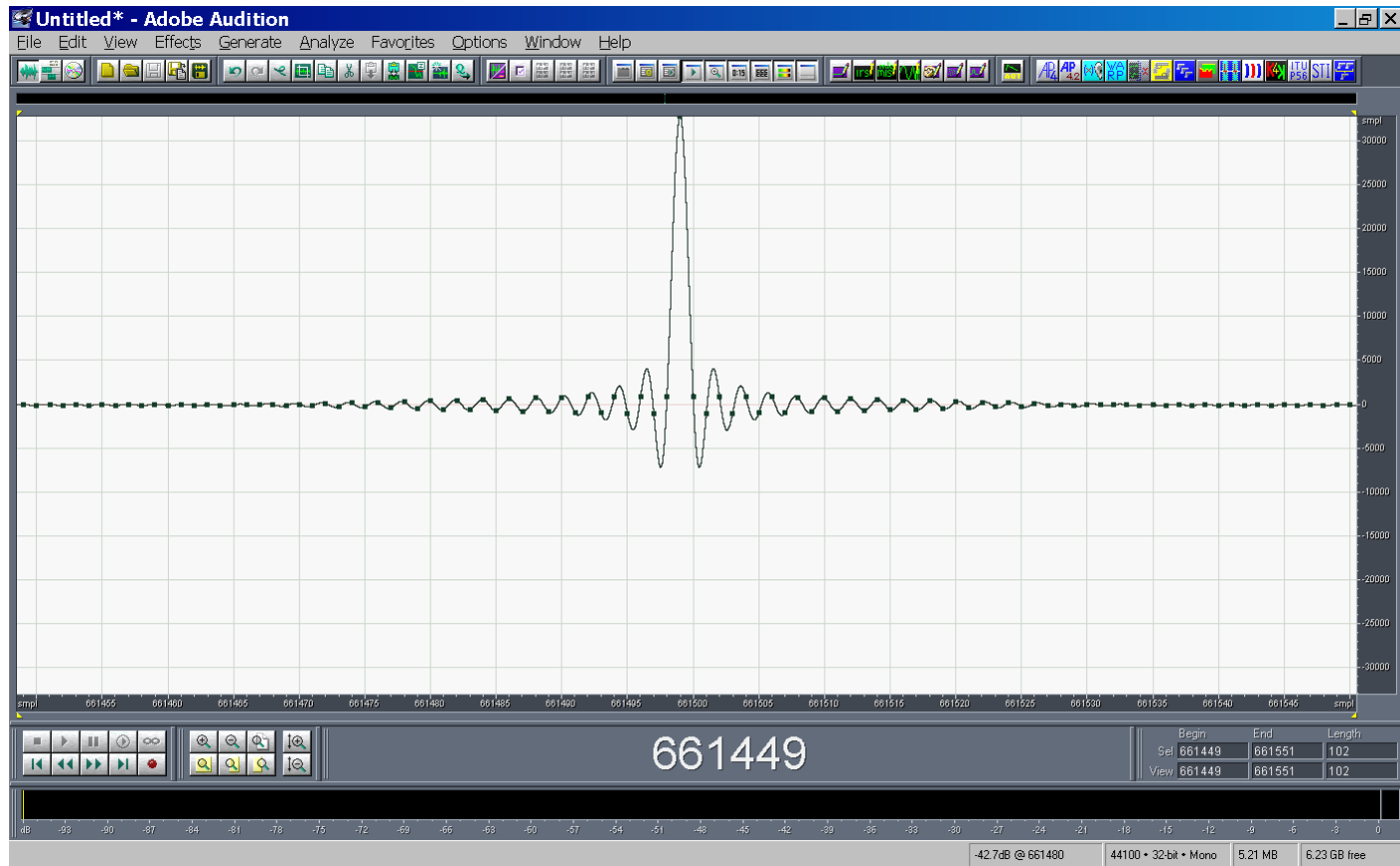
# Topics

- Pre-ringing at high and low frequency before the arrival of the direct sound pulse
- Pre/post equalization of the test signal performed in a way which avoids time-smearing of the impulse response
- Sensitivity to abrupt pulsive noises during the measurement
- Skewing of the measured impulse response when the playback and recording digital clocks are mismatched
- Cancellation of high frequencies in the late part of the tail when performing synchronous averaging



# Pre-ringing at high and low frequency

- Pre-ringing at high frequency due to improper fade-out

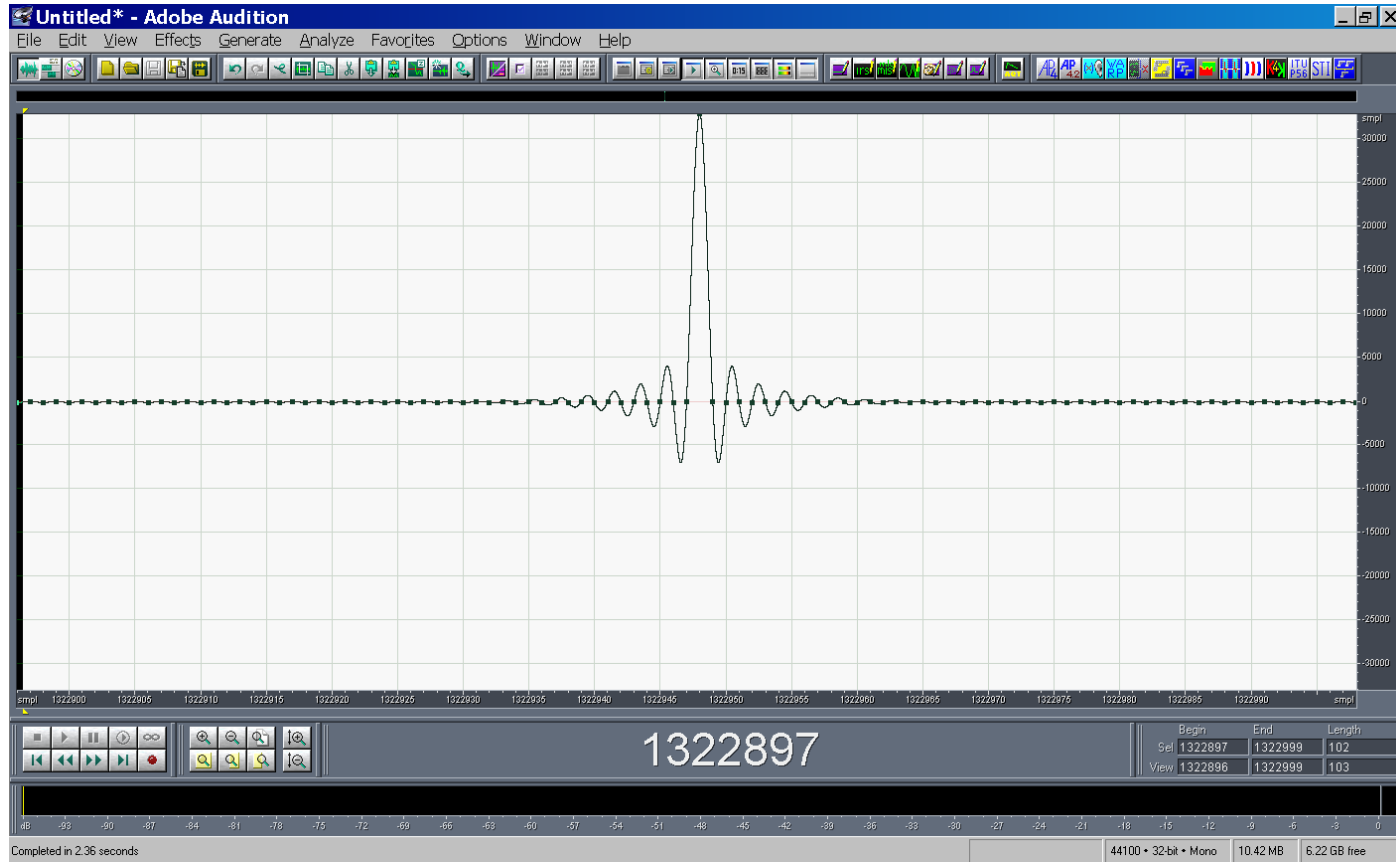


This picture shows the preringing obtained deconvolving directly the test signal, without passing through the system under test



# Pre-ringing at high and low frequency

- Perfect Dirac's delta after removing the fade-out



This picture shows the result obtained deconvolving directly the test signal, without passing through the system under test, and employing a sine sweep going up to the Nyquist frequency



# Pre-ringing at high and low frequency

- Pre-ringing at low frequency due to a bad sound card featuring frequency-dependent latency



This artifact can be corrected if the frequency-dependent latency remains the same, by creating a suitable inverse filter with the Kirkeby method





# Kirkeby inverse filter

- The Kirkeby inverse filter is computed inverting the measured IR

1) The IR to be inverted is FFT transformed to frequency domain:

$$H(f) = \text{FFT} [h(f)]$$

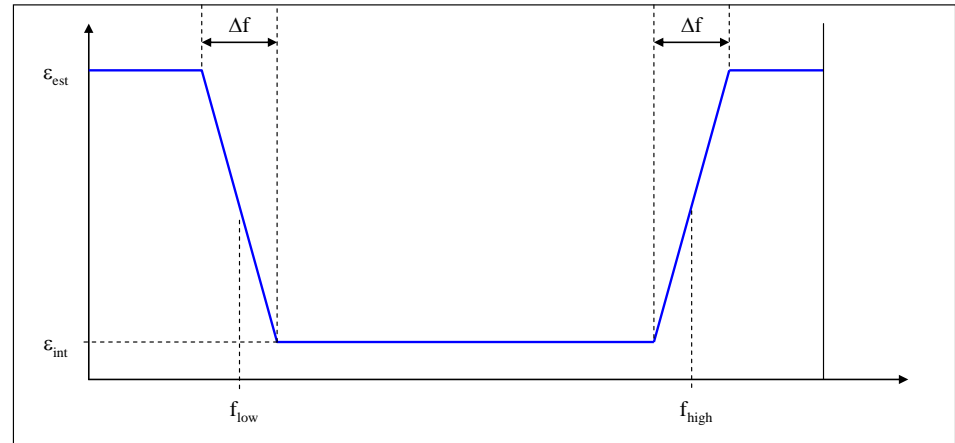
2) The computation of the inverse filter is done in frequency domain:

$$C(f) = \frac{\text{Conj}[H(f)]}{\text{Conj}[H(f)] \cdot H(f) + \varepsilon(f)}$$

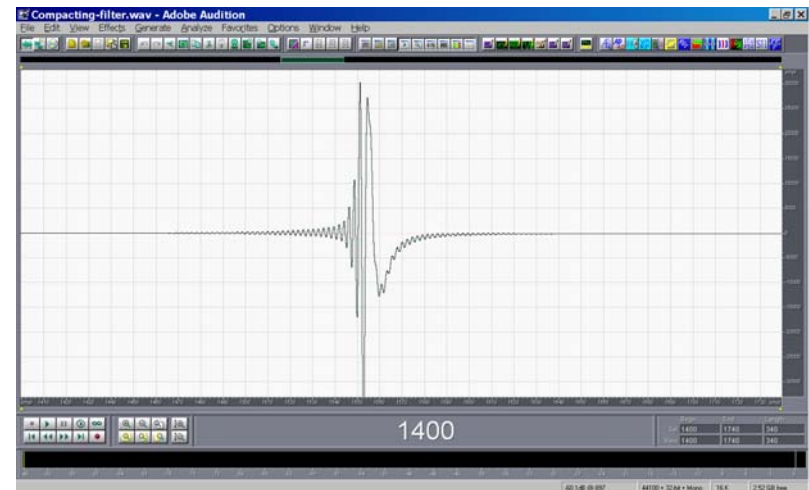
Where  $\varepsilon(f)$  is a small, frequency-dependent regularization parameter

3) Finally, an IFFT brings back the inverse filter to time domain:

$$c(t) = \text{IFFT} [C(f)]$$



Frequency-dependent regularization parameter

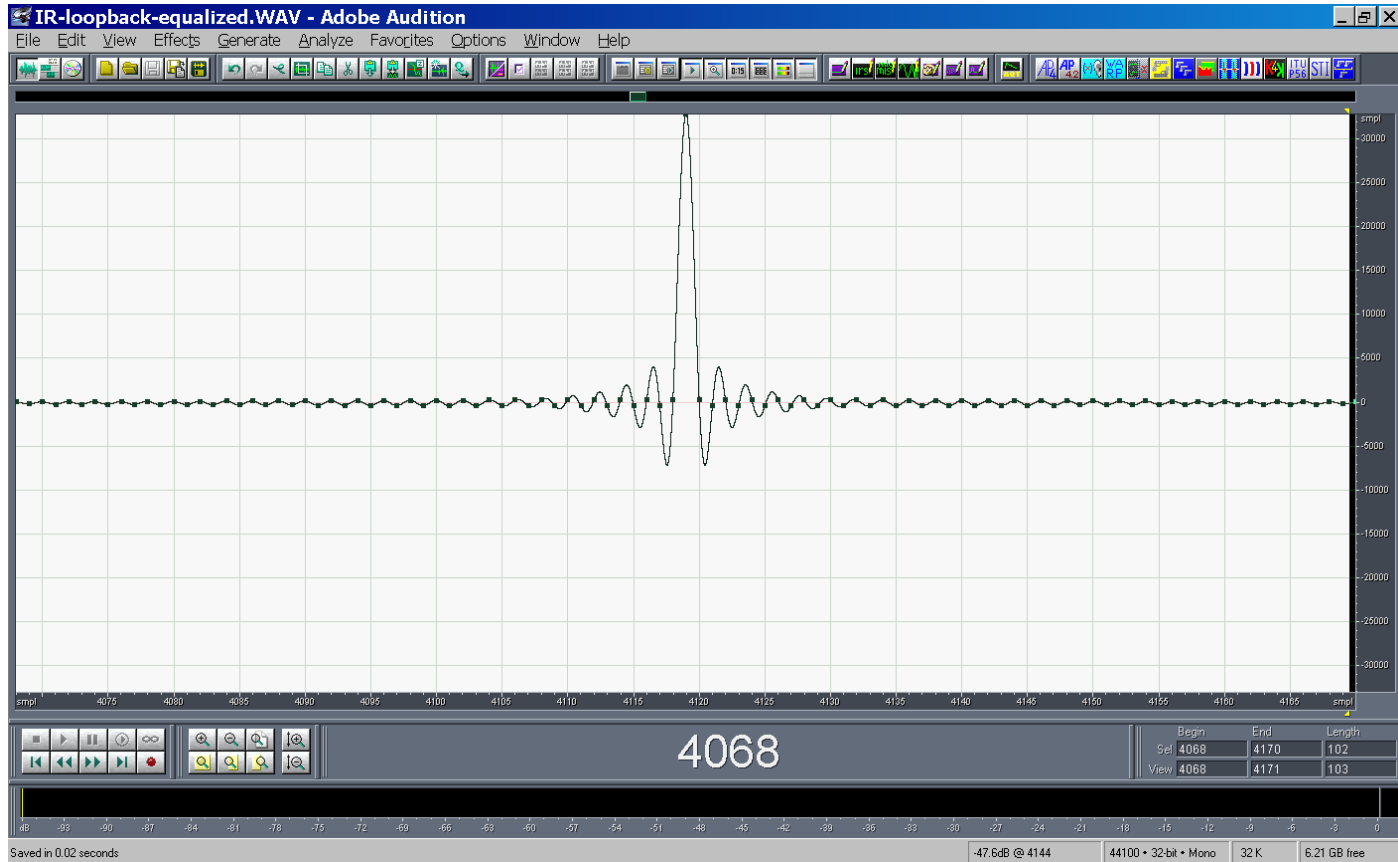


Inverse filter



# Pre-ringing at high and low frequency

- Convolution of the time-smeared IR with the Kirkeby compacting filter, a very sharp IR is obtained



The same method can also be applied for correcting the response of the loudspeaker/microphone system, if an anechoic preliminary test is done



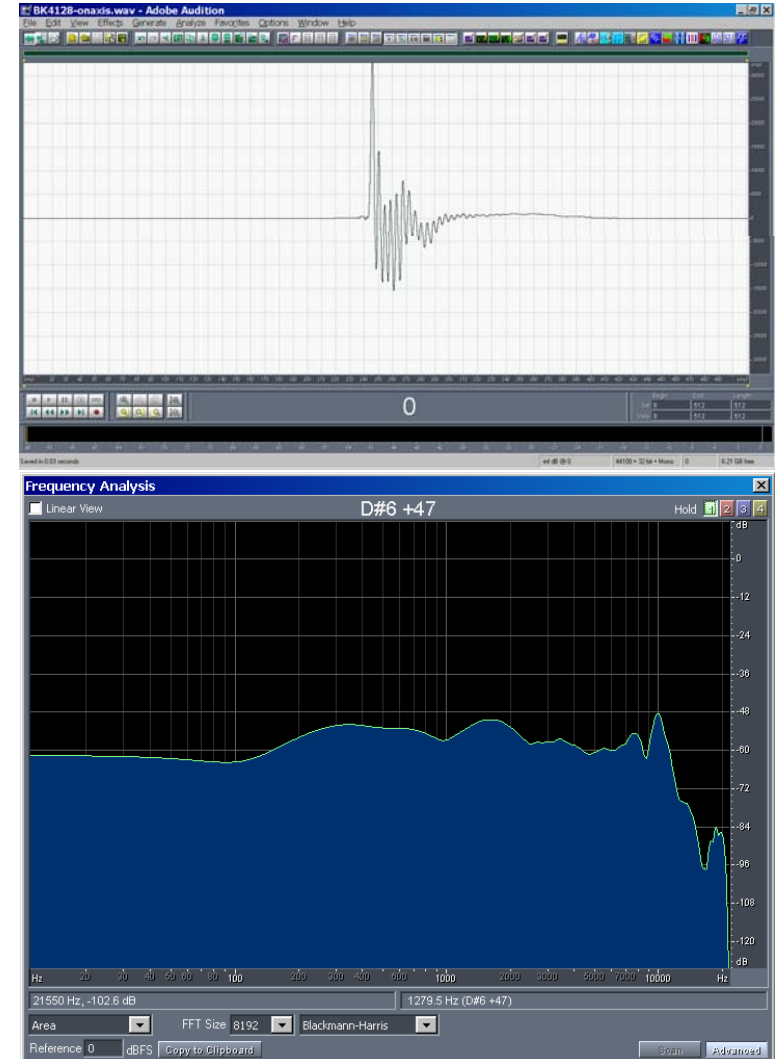
# Topics

- Pre-ringing at high and low frequency before the arrival of the direct sound pulse
- Pre/post equalization of the test signal performed in a way which avoids time-smearing of the impulse response
- Sensitivity to abrupt pulsive noises during the measurement
- Skewing of the measured impulse response when the playback and recording digital clocks are mismatched
- Cancellation of high frequencies in the late part of the tail when performing synchronous averaging



# Equalization of the whole system

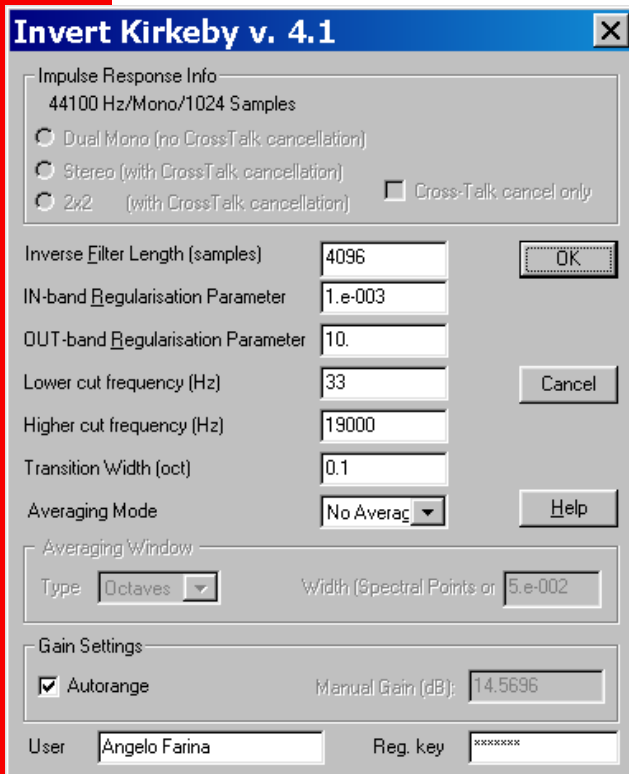
- An anechoic measurement is first performed





# Equalization of the whole system

- A suitable inverse filter is generated with the Kirkeby method by inverting the anechoic measurement

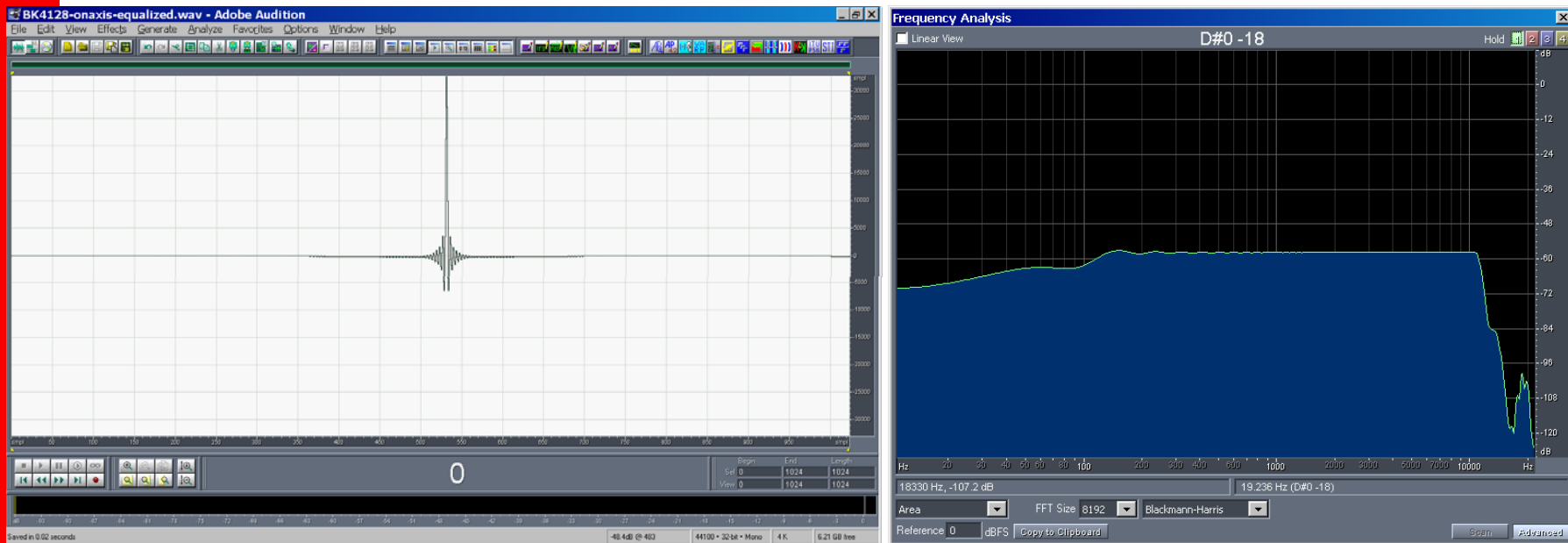






# Equalization of the whole system

- The inverse filter can be either pre-convolved with the test signal or post-convolved with the result of the measurement
- Pre-convolution usually reduces the SPL being generated by the loudspeaker, resulting in worst S/N ratio
- On the other hand, post-convolution can make the background noise to become “coloured”, and hence more perceptible
- The resulting anechoic IR becomes almost perfectly a Dirac’s Delta function:





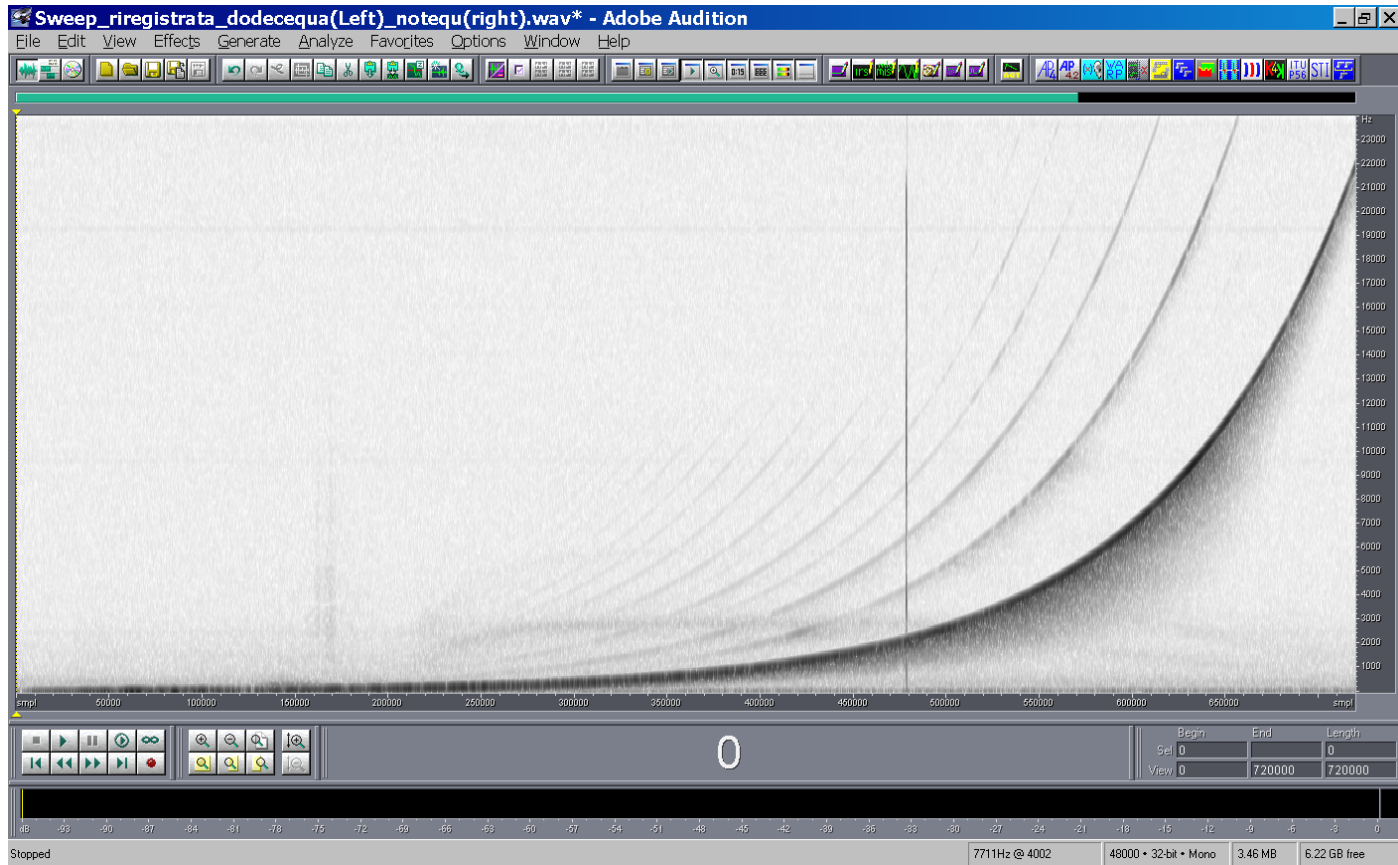
# Topics

- Pre-ringing at high and low frequency before the arrival of the direct sound pulse
- Pre/post equalization of the test signal performed in a way which avoids time-smearing of the impulse response
- Sensitivity to abrupt pulsive noises during the measurement
- Skewing of the measured impulse response when the playback and recording digital clocks are mismatched
- Cancellation of high frequencies in the late part of the tail when performing synchronous averaging



# Sensitivity to abrupt pulsive noises

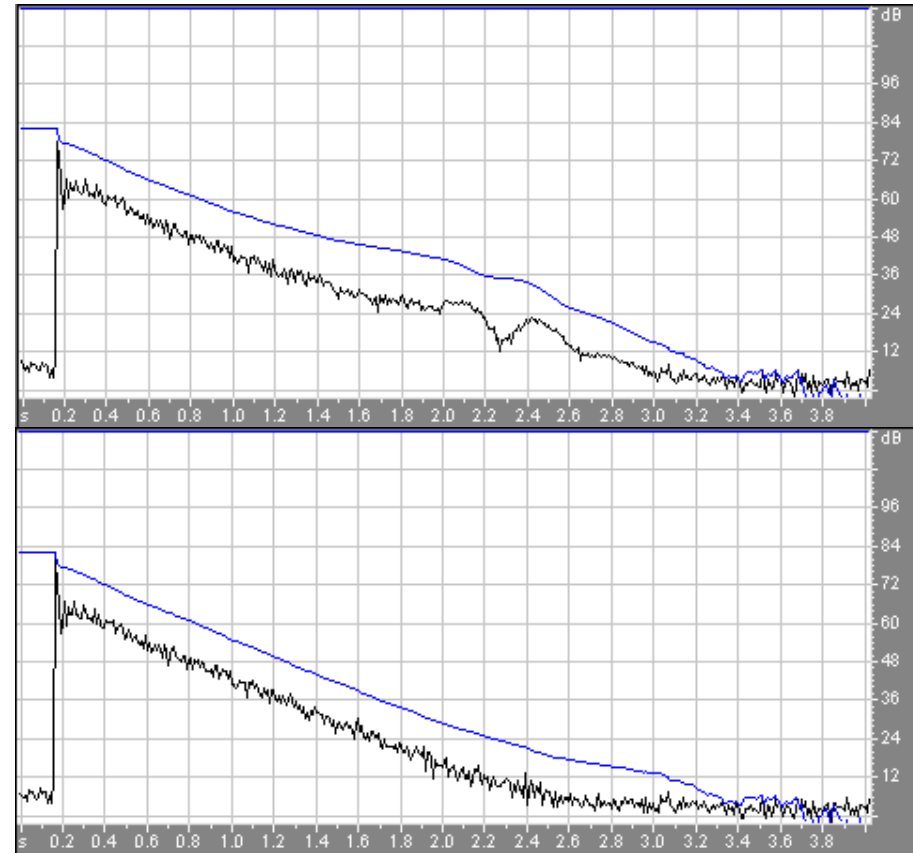
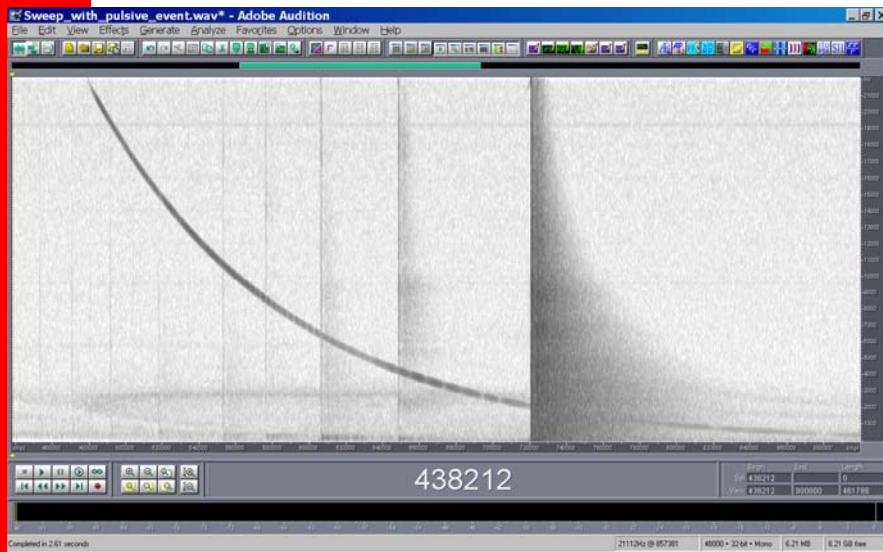
- Often a pulsive noise occurs during a sine sweep measurement





# Sensitivity to abrupt pulsive noises

- After deconvolution, the pulsive sound causes intolerable artifacts in the impulse response

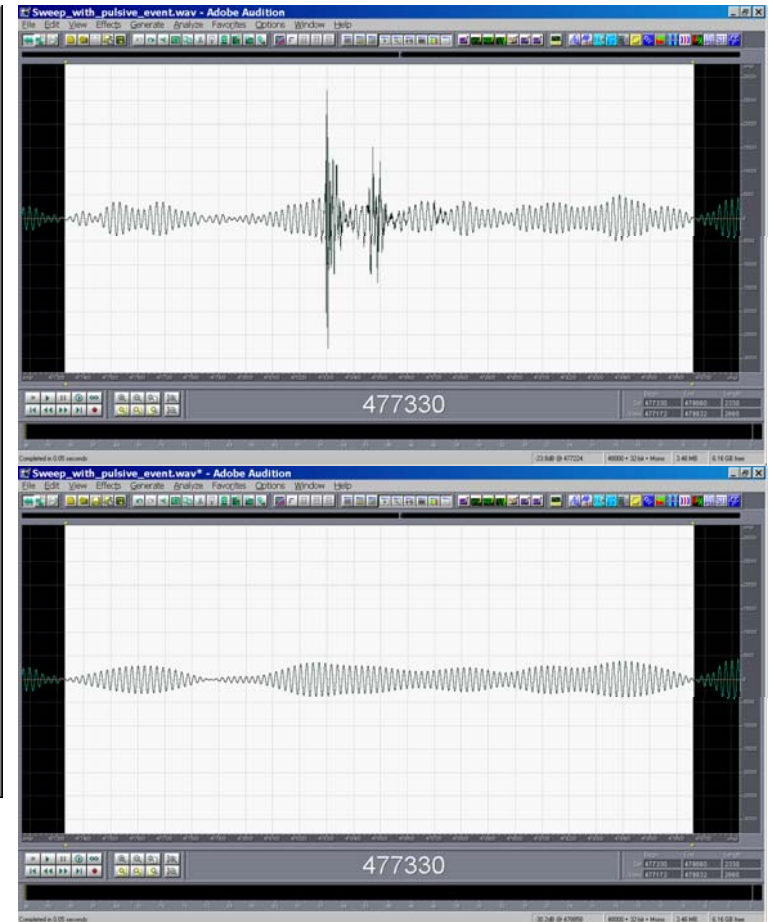
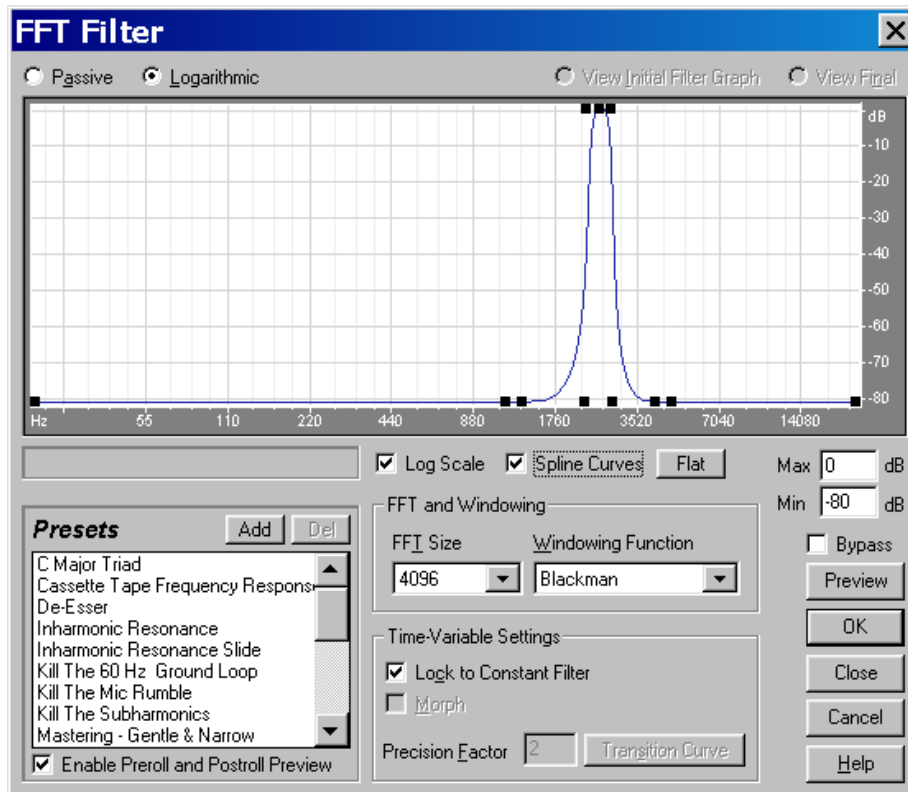


The artifact appears as a down-sloping sweep on the impulse response.  
At the 2 kHz octave band the decay is distorted, and the reverb. time is artificially increased from 2.13 to 2.48 s



# Sensitivity to abrupt pulsive noises

- **Several denoising techniques can be employed:**
  - ▶ Brutely silencing the transient noise
  - ▶ Employing the specific “click-pop eliminator” plugin of Adobe Audition
  - ▶ Applying a narrow-passband filter around the frequency which was being generated in the moment in which the pulsive noise occurred
- **The third approach provides the better results:**





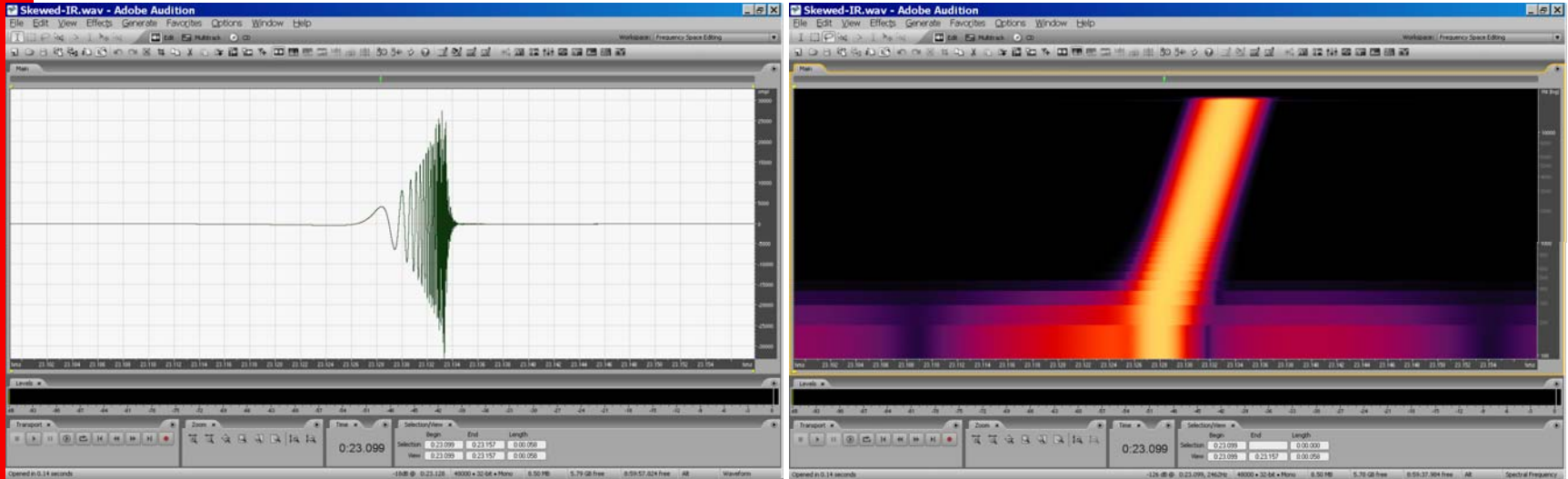


# Topics

- Pre-ringing at high and low frequency before the arrival of the direct sound pulse
- Pre/post equalization of the test signal performed in a way which avoids time-smearing of the impulse response
- Sensitivity to abrupt pulsive noises during the measurement
- Skewing of the measured impulse response when the playback and recording digital clocks are mismatched
- Cancellation of high frequencies in the late part of the tail when performing synchronous averaging

# Clock mismatch

- When the measurement is performed employing devices which exhibit significant clock mismatch between playback and recording, the resulting impulse response is “skewed” (stretched in time):

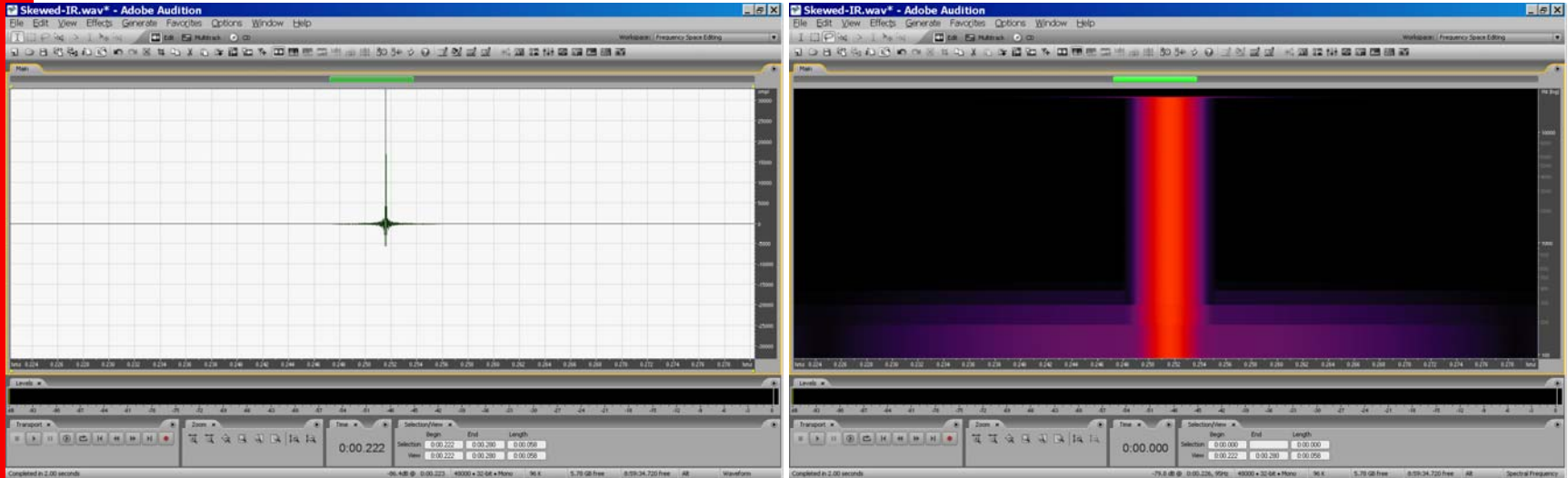


The pictures show the results of an electrical measurement performed connecting directly a CD-player with a DAT recorder



# Clock mismatch

- It is possible to re-pack the impulse response employing the already-described approach based on the usage of a Kirkeby inverse filter:

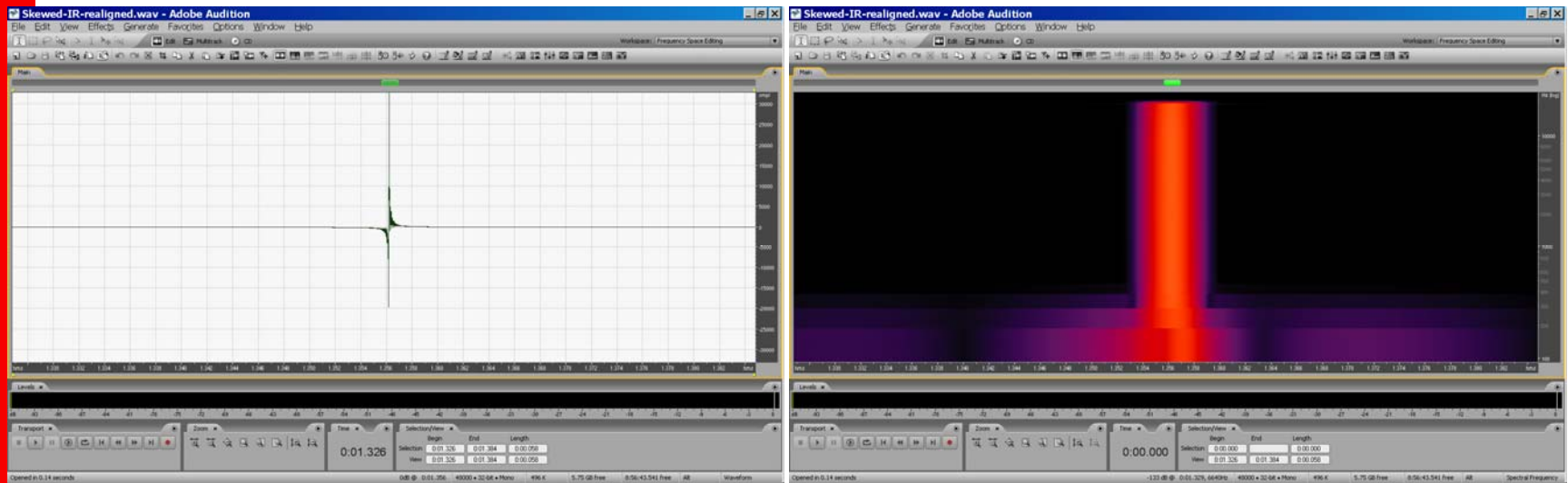


However, this is possible only if a “reference” electrical (or anechoic) measurement has been performed. But, in many cases, one only gets the re-recorded signals, and no reference measurement is available, so the Kirkeby inverse filter cannot be computed.



# Clock mismatch

- However, it is always possible to generate a pre-stretched inverse filter, which is longer or shorter than the “theoretical” one - by proper selection of the length of the inverse filter, it is possible to deconvolve impulse responses which are almost perfectly “unskewed”:



The pictures show the result of the deconvolution of a clock-mismatched measurement, in which a pre-stretched inverse filter is employed, 8.5 ms longer than the theoretical one.



# Topics

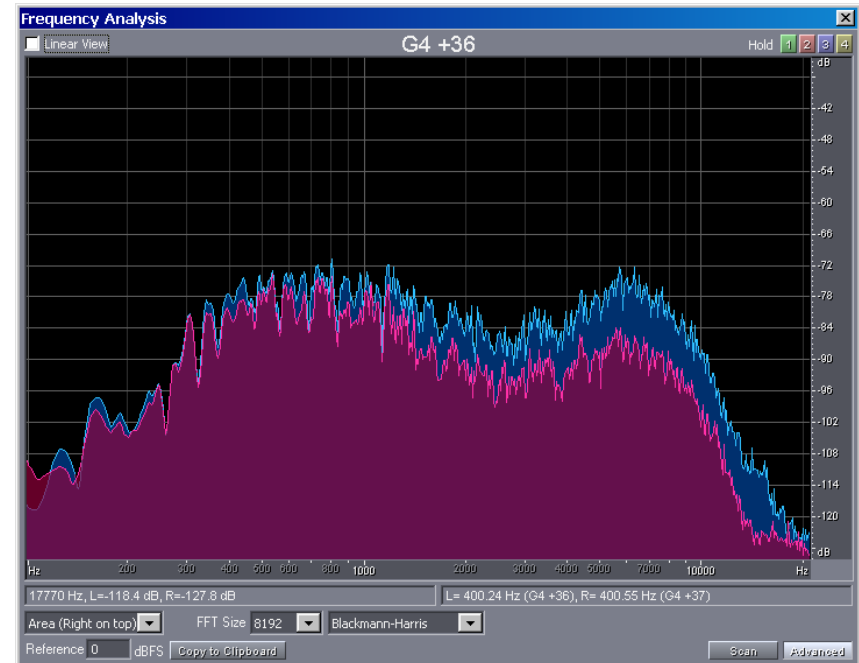
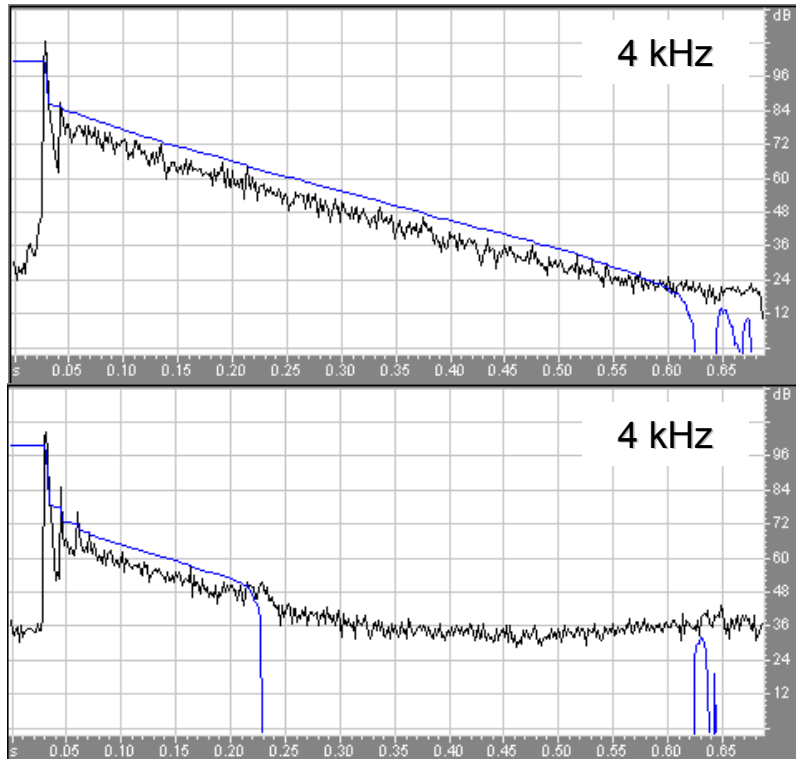
- Pre-ringing at high and low frequency before the arrival of the direct sound pulse
- Pre/post equalization of the test signal performed in a way which avoids time-smearing of the impulse response
- Sensitivity to abrupt pulsive noises during the measurement
- Skewing of the measured impulse response when the playback and recording digital clocks are mismatched
- Cancellation of high frequencies in the late part of the tail when performing synchronous averaging





# High-frequency cancellation due to averaging

- When several impulse response measurements are synchronously-averaged for improving the S/N ratio, the late part of the tail cancels out, particularly at high frequency, due to slight time variance of the system



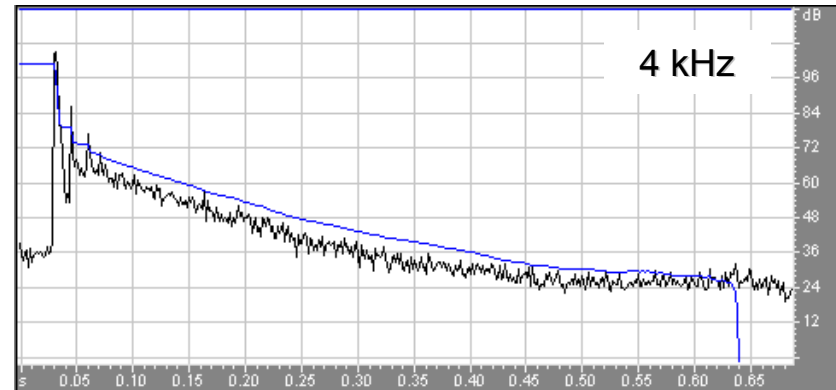
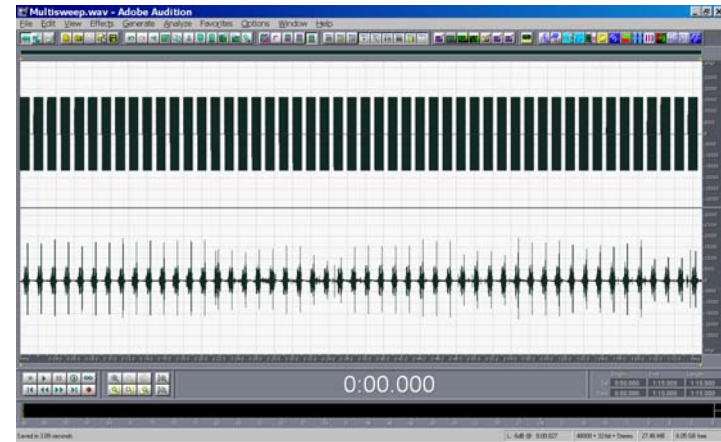
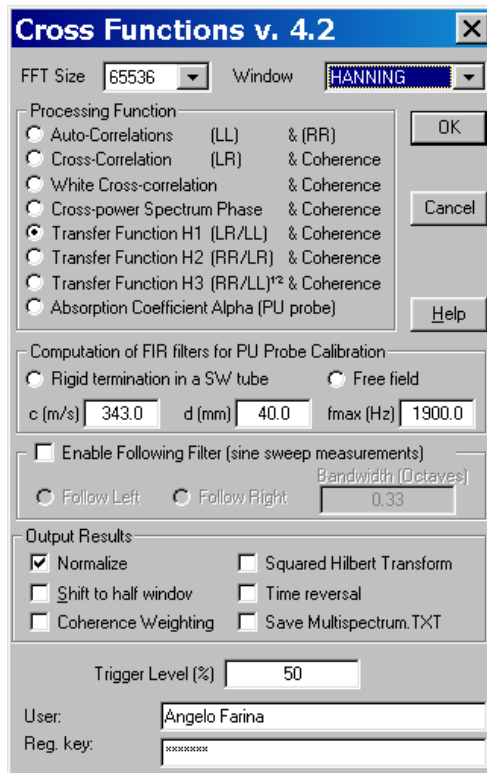
Spectrum of a single sweep of 50s (above) versus 50 sweeps of 1s (below) short-FFT spectrum at 200 ms after direct sound

Comparison of a single sweep 50 s long with the synchronous average of 50 sweeps, 1 s long each.



# High-frequency cancellation due to averaging

- However, if averaging is performed properly in spectral domain, and a single conversion to time domain is performed after averaging, this artifact is significantly reduced
- The new “cross Functions” plugin can be used for computing H1:



Result of transfer function H1, processing a sequence of 50 sine sweeps (above)

$$H_1(f) = \frac{G_{LR}}{G_{LL}}$$



# Conclusions (I of II)

- **The ESS method revealed to be systematically superior to the MLS method for measuring electroacoustical impulse responses**
- **Traditional methods for measuring “spatial parameters” (IACC, LF) proved to be unreliable and do not provide complete information**
- **The 1°-order Ambisonics method can be used for generating and recording sound with a limited amount of spatial information**
- **For obtaining better spatial resolution, High-Order Ambisonics can be used, limiting the spherical-harmonics expansion to a reasonable order (2°, 3° or 4°).**
- **Experimental hardware and software tools have been developed (mainly in France, but also in USA), allowing to build an inexpensive complete measurement system**



# Conclusions (II of II)

- **ESS is now employed in top-grade measurement systems: Audio Precision (TM), Rhode-Schwartz and B&K / DIRAC**
- **However, these completely-packaged measurement systems often do not allow to play “tricks” and to adjust the signals for solving problems, which have been shown here**
- **Workarounds have been found for the problems occurring when performing ESS measurements**
- **These workarounds are easily applied by working with a general purpose sound editor (Adobe Audition)**
- **A number of additional plugins have been developed, making easy to generate the test signal, to deconvolve and process impulse responses, to compute inverse filters and to perform advanced processing (STI, AQT, etc.)**
- **These plugins are freely downloadable at the AURORA web site:**

**[www.aurora-plugins.com](http://www.aurora-plugins.com)**